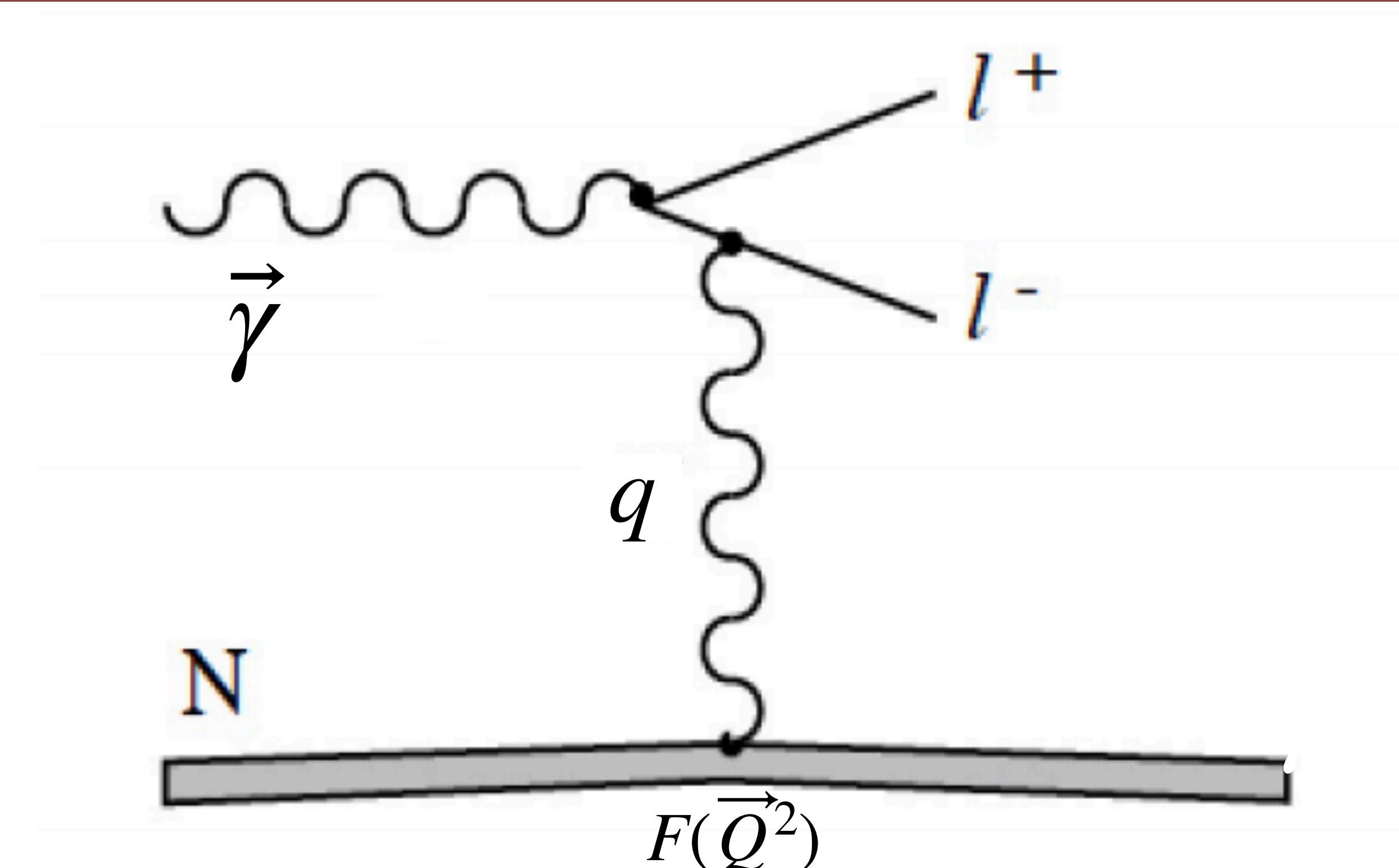


Using Bethe Heitler Pairs as a Polarimeter in GlueX



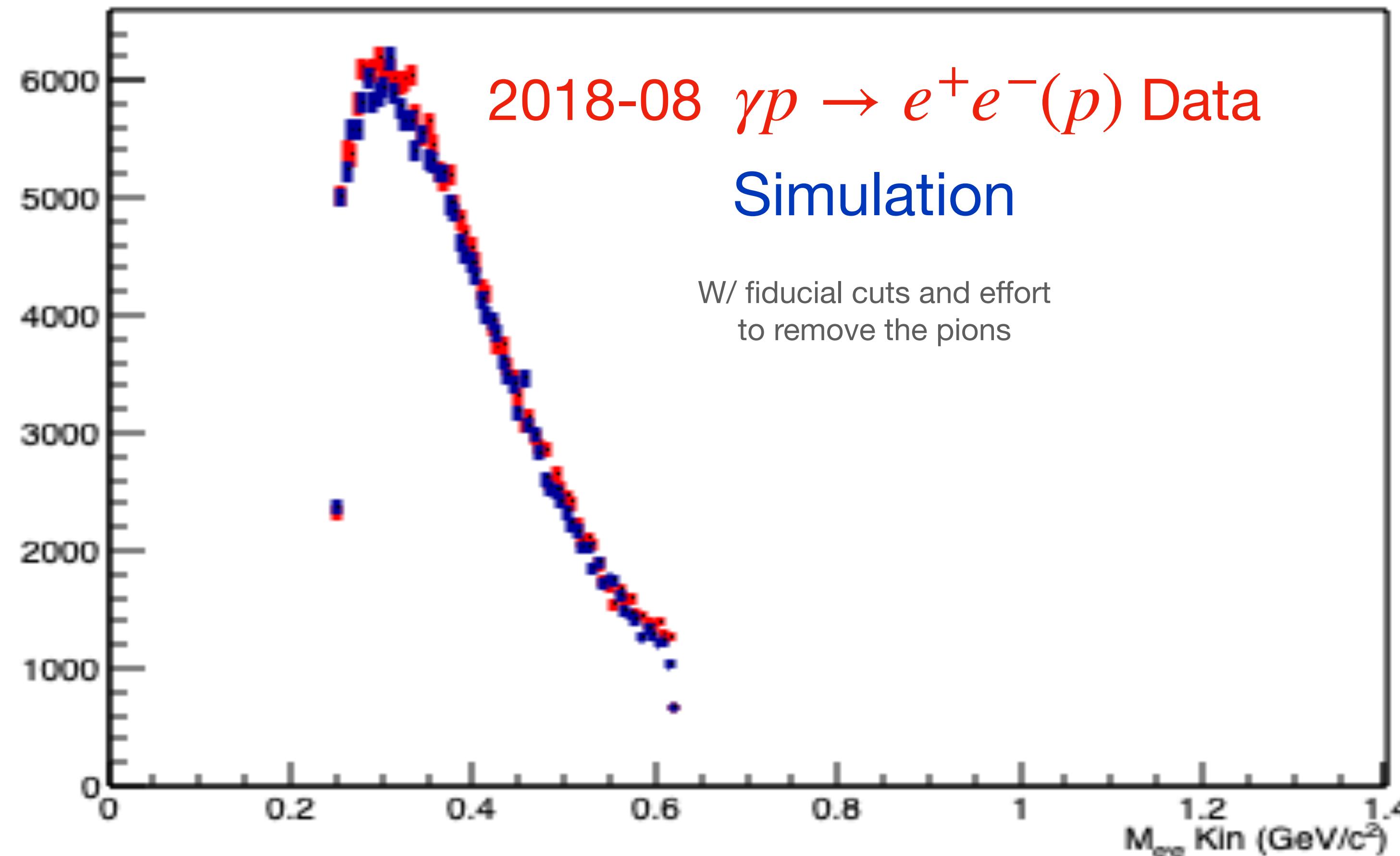
Andrew Schick

GlueX Spring Collaboration Meeting, Friday, May 28 2021

Motivation

- Have a method of verifying the result from TPOL in the CPP experiment
- **Advantages**
 - For sufficiently long runs you get this method “for free”
 - High analyzing power ~60%
- **Challenges**
 - Addressing/removing pion background

e^+e^- Invariant Mass



**Neural Net to sort pions
and electrons**

**Fiducial cuts to get agreement
between data and
simulation**

Still actively studying

Use Bethe-Heitler pairs to measure linear photon polarization.

$$\begin{aligned} d\sigma &= \left(\frac{1 + \mathcal{P}}{2} \right) d\sigma_{||} + \left(\frac{1 - \mathcal{P}}{2} \right) d\sigma_{\perp} \\ &= \left(\frac{d\sigma_{||} + d\sigma_{\perp}}{2} \right) + \mathcal{P} \left(\frac{d\sigma_{||} - d\sigma_{\perp}}{2} \right) \end{aligned}$$



Use Bethe-Heitler pairs to measure linear photon polarization.

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$$\uparrow \\ d\sigma_0$$

Unpolarized

$$\uparrow \\ d\sigma_1$$

Polarized

Bakmaev et al, Physics Letters B 660 (2008) 494-500
Modern Vectorized Approach

$$\vec{J}_T = \frac{\vec{p}_1}{p_1^2 + m^2} + \frac{\vec{p}_2}{p_2^2 + m^2} = \frac{\vec{p}_1}{c_1} + \frac{\vec{p}_2}{c_2}$$

\vec{p}_1, \vec{p}_2 are the lepton's transverse momenta

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$d\sigma_0$ $d\sigma_1$

Unpolarized **Polarized**

Bakmaev et al, Physics Letters B 660 (2008) 494-500

Modern Vectorized Approach

$$\vec{J}_T = \frac{\vec{p}_1}{p_1^2 + m^2} + \frac{\vec{p}_2}{p_2^2 + m^2} = \frac{\vec{p}_1}{c_1} + \frac{\vec{p}_2}{c_2}$$

\vec{p}_1 , \vec{p}_2 are the lepton's transverse momenta

$$d\sigma_1 \sim P_\gamma |\vec{J}_T|^2 \cos(2\phi_{J_T})$$

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$$\begin{aligned} d\sigma &= \left(\frac{1 + \mathcal{P}}{2} \right) d\sigma_{||} + \left(\frac{1 - \mathcal{P}}{2} \right) d\sigma_{\perp} \\ &= \left(\frac{d\sigma_{||} + d\sigma_{\perp}}{2} \right) + \mathcal{P} \left(\frac{d\sigma_{||} - d\sigma_{\perp}}{2} \right) \end{aligned}$$

$d\sigma_0$

Unpolarized

$$d\sigma_1$$


Polarized

Bakmaev et al, Physics Letters B 660 (2008) 494-500 Modern Vectorized Approach

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\vec{p}_1 , \vec{p}_2 are the lepton's transverse momenta

$$d\sigma_1 \sim P_\gamma |\vec{J}_T|^2 \cos(2\phi_{J_T})$$

Bakmaev's formulation is really only valid at very large t

Vectorizing the Classic Bethe-Heitler Formulation

$$d\sigma = \frac{Z^2}{137} \frac{e^4}{4\pi^2} \frac{\not{p}_+ \not{p}_- dE_+ d\Omega_+ d\Omega_-}{k^3 q^4} \left\{ \frac{(\boldsymbol{\epsilon} \cdot \mathbf{p}_+)^2 (q^2 - 4E_-^2)}{(E_+ - \not{p}_+ \cos\theta_+)^2} \right. \\ + \frac{(\boldsymbol{\epsilon} \cdot \mathbf{p}_-)^2 (q^2 - 4E_+^2)}{(E_- - \not{p}_- \cos\theta_-)^2} - \frac{2(\boldsymbol{\epsilon} \cdot \mathbf{p}_+) (\boldsymbol{\epsilon} \cdot \mathbf{p}_-) (q^2 + 4E_+ E_-)}{(E_+ - \not{p}_+ \cos\theta_+) (E_- - \not{p}_- \cos\theta_-)} \\ \left. + \frac{k^2 [\not{p}_+^2 \sin^2\theta_+ + \not{p}_-^2 \sin^2\theta_- + 2\not{p}_+ \not{p}_- \sin\theta_+ \sin\theta_- \cos(\varphi_\tau - \varphi_-)]}{(E_+ - \not{p}_+ \cos\theta_+) (E_- - \not{p}_- \cos\theta_-)} \right\}.$$

T.H. Berlin and L. Madansky, Phys. Rev. **78**, 623 (1950)

$\boldsymbol{\epsilon}$ is a unit vector in the direction of polarization of the incident photon.

Vectorizing the Classic Bethe-Heitler Formulation

$$\vec{J}_T = \frac{2E_2}{E_1 - p_1 \cos \theta_1} \vec{p}_1 T + \frac{2E_2}{E_2 - p_2 \cos \theta_2} \vec{p}_2 T$$

$$\vec{K}_T = \frac{\sqrt{q^2}}{E_1 - p_1 \cos \theta_1} \vec{p}_1 T - \frac{\sqrt{q^2}}{E_2 - p_2 \cos \theta_2} \vec{p}_2 T$$

$$d\sigma = \frac{Z^2}{137} \frac{e^4}{4\pi^2} \frac{p_+ p_- dE_+ d\Omega_+ d\Omega_-}{k^3 q^4} \left\{ \frac{(\boldsymbol{\epsilon} \cdot \mathbf{p}_+)^2 (q^2 - 4E_-^2)}{(E_+ - p_+ \cos \theta_+)^2} \right. \\ \left. + \frac{(\boldsymbol{\epsilon} \cdot \mathbf{p}_-)^2 (q^2 - 4E_+^2)}{(E_- - p_- \cos \theta_-)^2} - \frac{2(\boldsymbol{\epsilon} \cdot \mathbf{p}_+) (\boldsymbol{\epsilon} \cdot \mathbf{p}_-) (q^2 + 4E_+ E_-)}{(E_+ - p_+ \cos \theta_+) (E_- - p_- \cos \theta_-)} \right. \\ \left. + \frac{k^2 [p_+^2 \sin^2 \theta_+ + p_-^2 \sin^2 \theta_- + 2p_+ p_- \sin \theta_+ \sin \theta_- \cos(\varphi_+ - \varphi_-)]}{(E_+ - p_+ \cos \theta_+) (E_- - p_- \cos \theta_-)} \right\}.$$

$\boldsymbol{\epsilon}$ is a unit vector in the direction of polarization of the incident photon.

Then:

Vectorizing the Classic Bethe-Heitler Formulation

$$\vec{J}_T = \frac{2E_2}{E_1 - p_1 \cos \theta_1} \vec{p}_{1T} + \frac{2E_2}{E_2 - p_2 \cos \theta_2} \vec{p}_{2T}$$

$$\vec{K}_T = \frac{\sqrt{q^2}}{E_1 - p_1 \cos \theta_1} \vec{p}_{1T} - \frac{\sqrt{q^2}}{E_2 - p_2 \cos \theta_2} \vec{p}_{2T}$$

Then:

$$d\sigma = d\sigma_0 + P_\gamma d\sigma_1$$

$$d\sigma_0 = \frac{d\sigma_{||} + d\sigma_{\perp}}{2} = k \left[- \left| \vec{J}_T \right|^2 + \left| \vec{K}_T \right|^2 + 2E_0^2 \frac{\left| \vec{p}_{1T} + \vec{p}_{2T} \right|^2}{(E_1 - p_1 \cos \theta_1)(E_2 - p_2 \cos \theta_2)} \right]$$

$$d\sigma_1 = \frac{d\sigma_{||} - d\sigma_{\perp}}{2} = k \left[- \left| \vec{J}_T \right|^2 \cos 2\phi_{J_T} + \left| \vec{K}_T \right|^2 \cos 2\phi_{J_T} \right]$$

$$d\sigma = \frac{Z^2}{137} \frac{e^4}{4\pi^2} \frac{p_+ p_- dE_+ d\Omega_+ d\Omega_-}{k^3 q^4} \left\{ \frac{(\boldsymbol{\epsilon} \cdot \mathbf{p}_+)^2 (q^2 - 4E_-^2)}{(E_+ - p_+ \cos \theta_+)^2} \right. \\ \left. + \frac{(\boldsymbol{\epsilon} \cdot \mathbf{p}_-)^2 (q^2 - 4E_+^2)}{(E_- - p_- \cos \theta_-)^2} - \frac{2(\boldsymbol{\epsilon} \cdot \mathbf{p}_+) (\boldsymbol{\epsilon} \cdot \mathbf{p}_-) (q^2 + 4E_+ E_-)}{(E_+ - p_+ \cos \theta_+) (E_- - p_- \cos \theta_-)} \right. \\ \left. + \frac{k^2 [p_+^2 \sin^2 \theta_+ + p_-^2 \sin^2 \theta_- + 2p_+ p_- \sin \theta_+ \sin \theta_- \cos(\varphi_+ - \varphi_-)]}{(E_+ - p_+ \cos \theta_+) (E_- - p_- \cos \theta_-)} \right\}.$$

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Then:

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$$d\sigma_0 = \frac{d\sigma_{||} + d\sigma_{\perp}}{2} = k \left[- \left| \vec{J}_T \right|^2 + \left| \vec{K}_T \right|^2 + 2E_0^2 \frac{\left| \vec{p}_{1T} + \vec{p}_{2T} \right|^2}{(E_1 - p_1 \cos \theta_1)(E_2 - p_2 \cos \theta_2)} \right]$$

$$d\sigma_1 = \frac{d\sigma_{||} - d\sigma_{\perp}}{2} = k \left[- \left| \vec{J}_T \right|^2 \cos 2\phi_{J_T} + \left| \vec{K}_T \right|^2 \cos 2\phi_{J_T} \right] \quad \left| \vec{J}_T \right|^2 \gg \left| \vec{K}_T \right|^2$$

$$d\sigma = \frac{Z^2}{137} \frac{e^4}{4\pi^2} \frac{p_+ p_- dE_+ d\Omega_+ d\Omega_-}{k^3 q^4} \left\{ \frac{(\boldsymbol{\epsilon} \cdot \mathbf{p}_+)^2 (q^2 - 4E_-^2)}{(E_+ - p_+ \cos \theta_+)^2} \right. \\ \left. + \frac{(\boldsymbol{\epsilon} \cdot \mathbf{p}_-)^2 (q^2 - 4E_+^2)}{(E_- - p_- \cos \theta_-)^2} - \frac{2(\boldsymbol{\epsilon} \cdot \mathbf{p}_+)(\boldsymbol{\epsilon} \cdot \mathbf{p}_-)(q^2 + 4E_+ E_-)}{(E_+ - p_+ \cos \theta_+)(E_- - p_- \cos \theta_-)} \right. \\ \left. + \frac{k^2 [p_+^2 \sin^2 \theta_+ + p_-^2 \sin^2 \theta_- + 2p_+ p_- \sin \theta_+ \sin \theta_- \cos(\varphi_+ - \varphi_-)]}{(E_+ - p_+ \cos \theta_+)(E_- - p_- \cos \theta_-)} \right\}.$$

$\boldsymbol{\epsilon}$ is a unit vector in the direction of polarization of the incident photon.

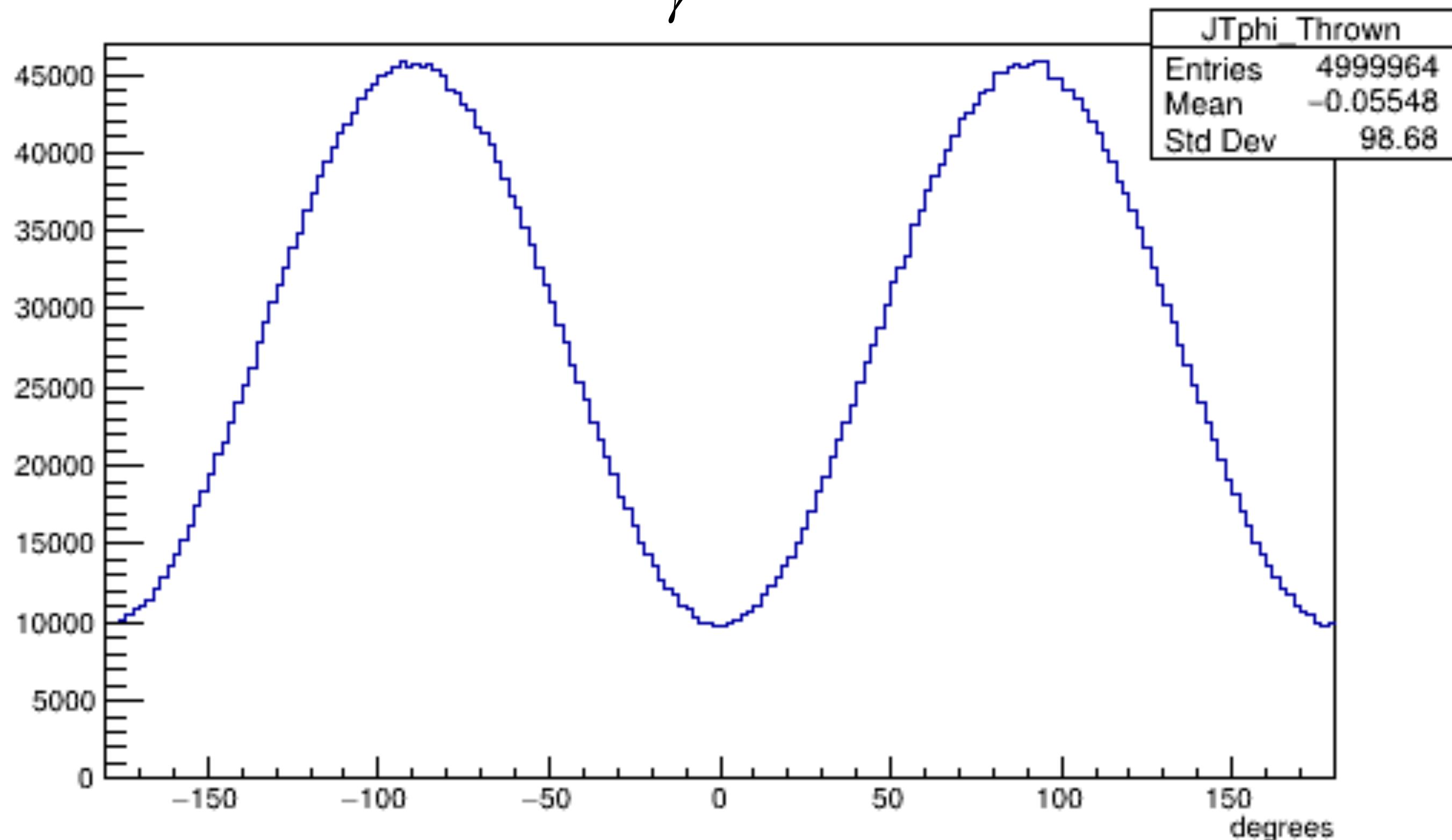
Plotting ϕ of \vec{J}_T from Monte Carlo

$$d\sigma = d\sigma_0 + P_\gamma d\sigma_1$$

$$d\sigma_1 = \sim \left| \vec{J}_T \right|^2 \cos 2\phi_{J_T}$$

MC with BH Cross-Section

$$P_\gamma = 1$$



1. Generate e+e- 4 vectors using this cross section

2. Plot ϕ_{J_T} from the 4 vectors

3. Measuring ϕ_{J_T} allows you to infer the beam polarization

2018-01 GlueX data

$\gamma p \rightarrow e^+e^- (p)$ Reaction Filter

Neural Net Cuts:

Neural Net Classification Cuts (NN1, NN2 > 0.8)

Fiducial Cuts:

$8.2 \text{ GeV} < E_\gamma < 8.8 \text{ GeV}$

$0.25 \text{ GeV} < W_{ee} < 0.621 \text{ GeV}$

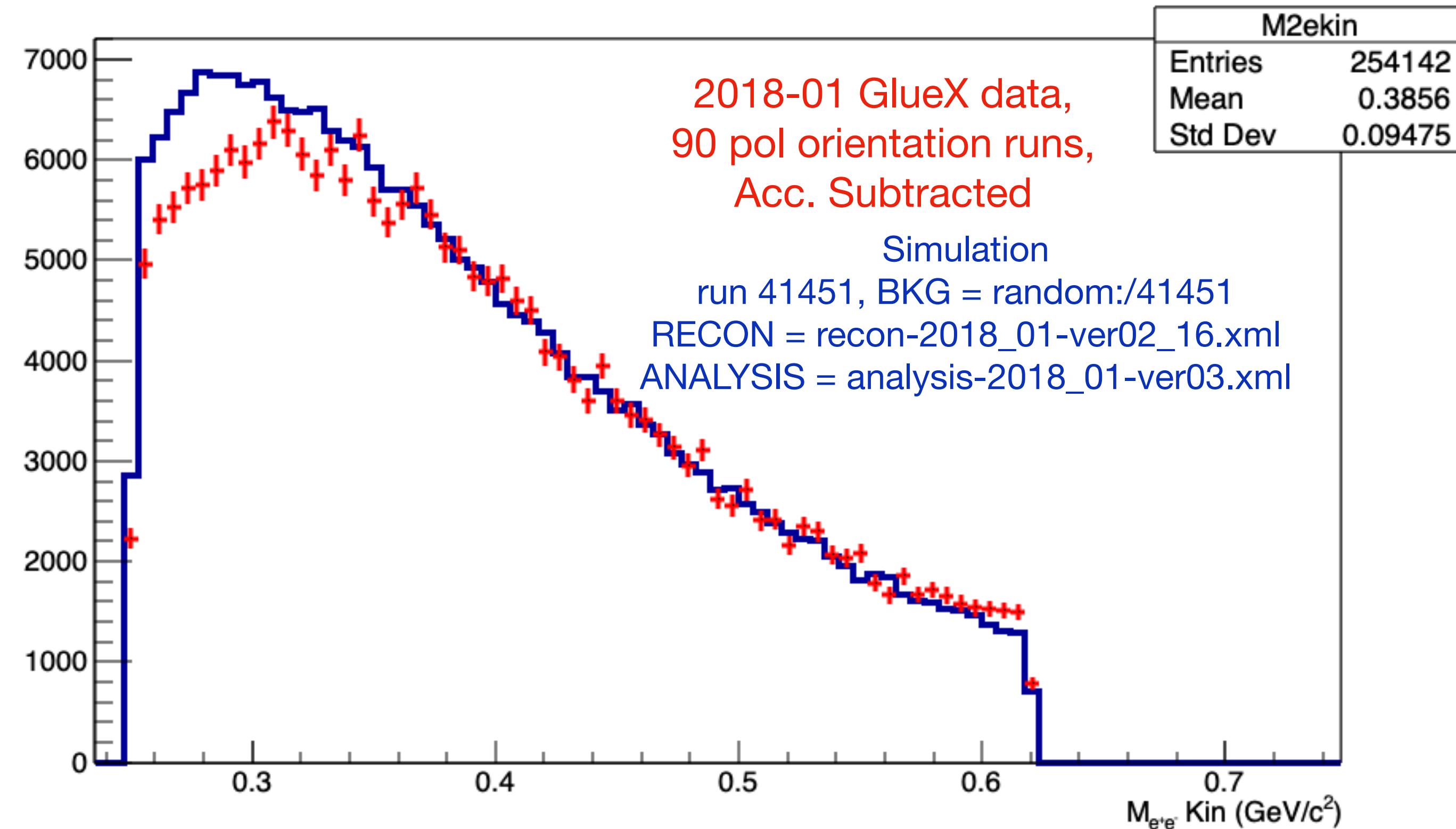
Both tracks have hits in the TOF

$\theta_1, \theta_2 > 1.5 \text{ deg}$

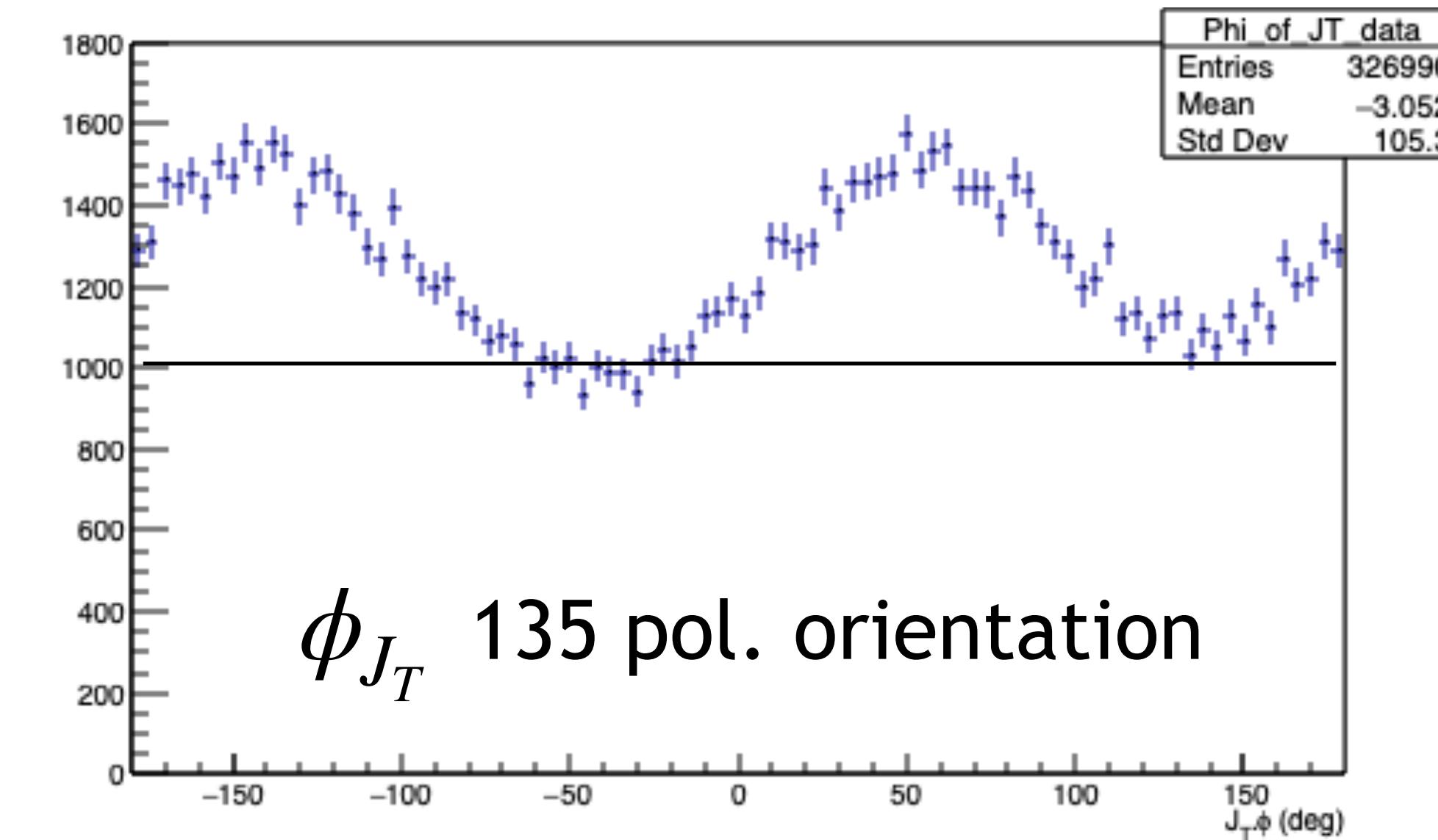
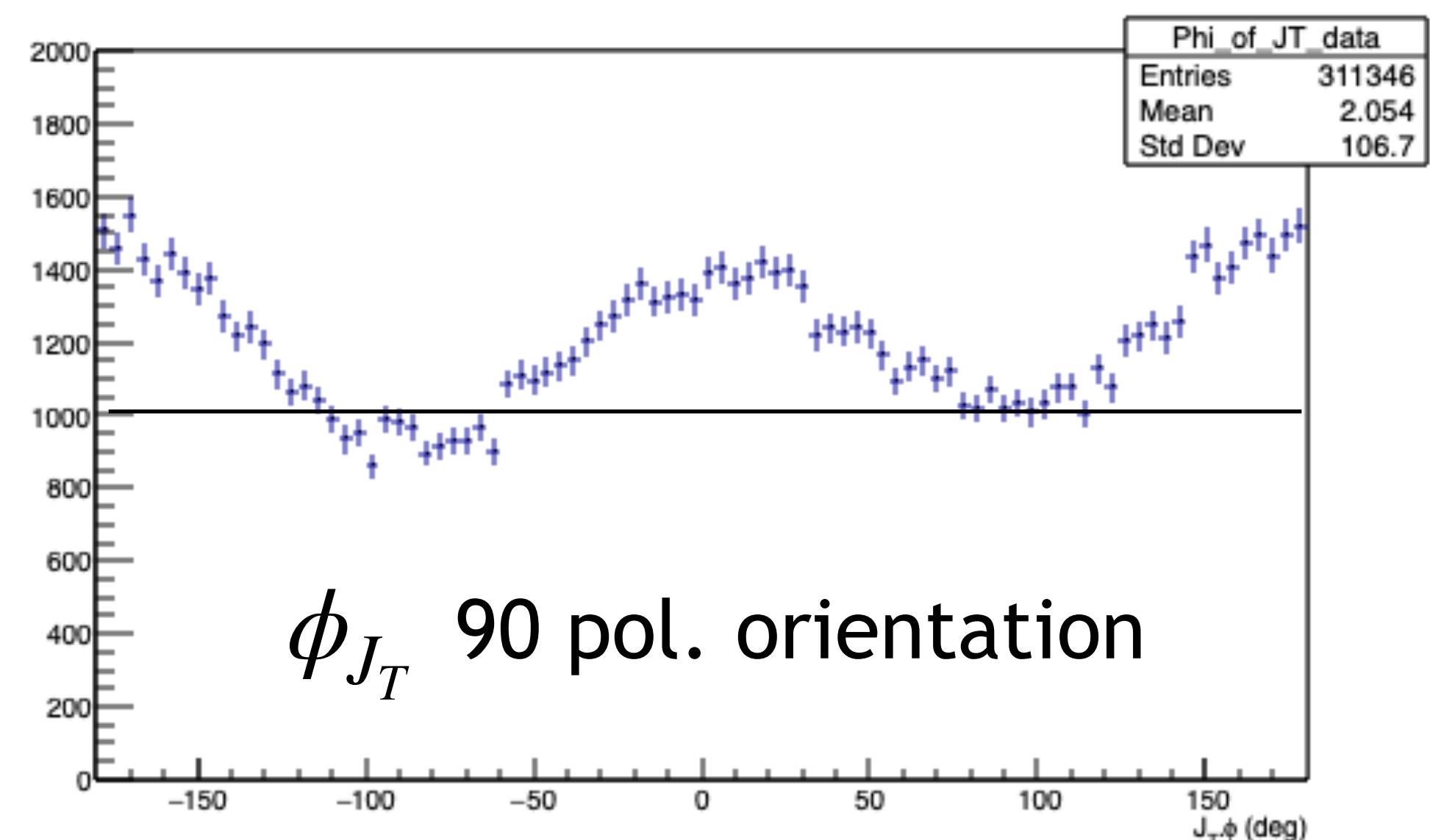
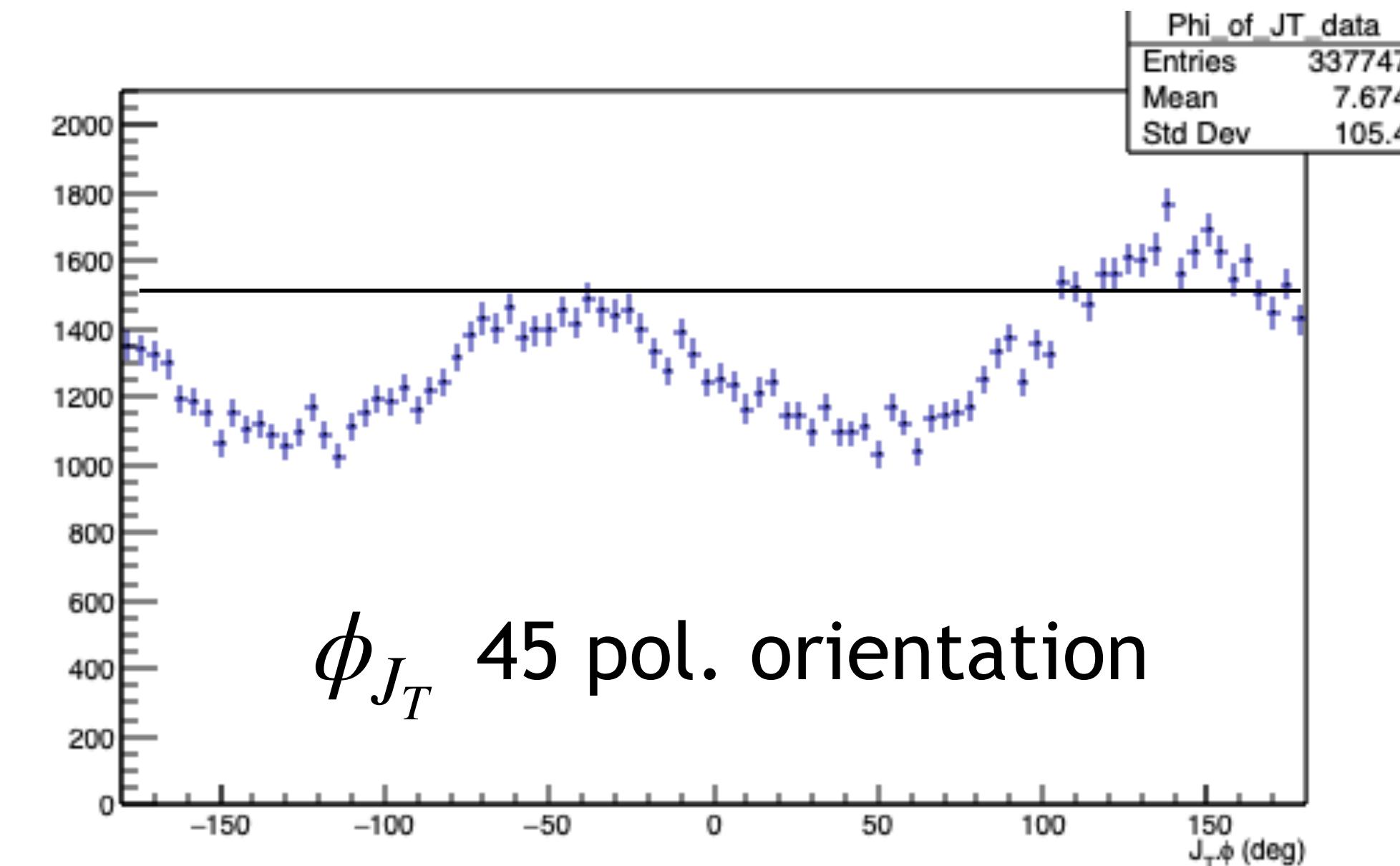
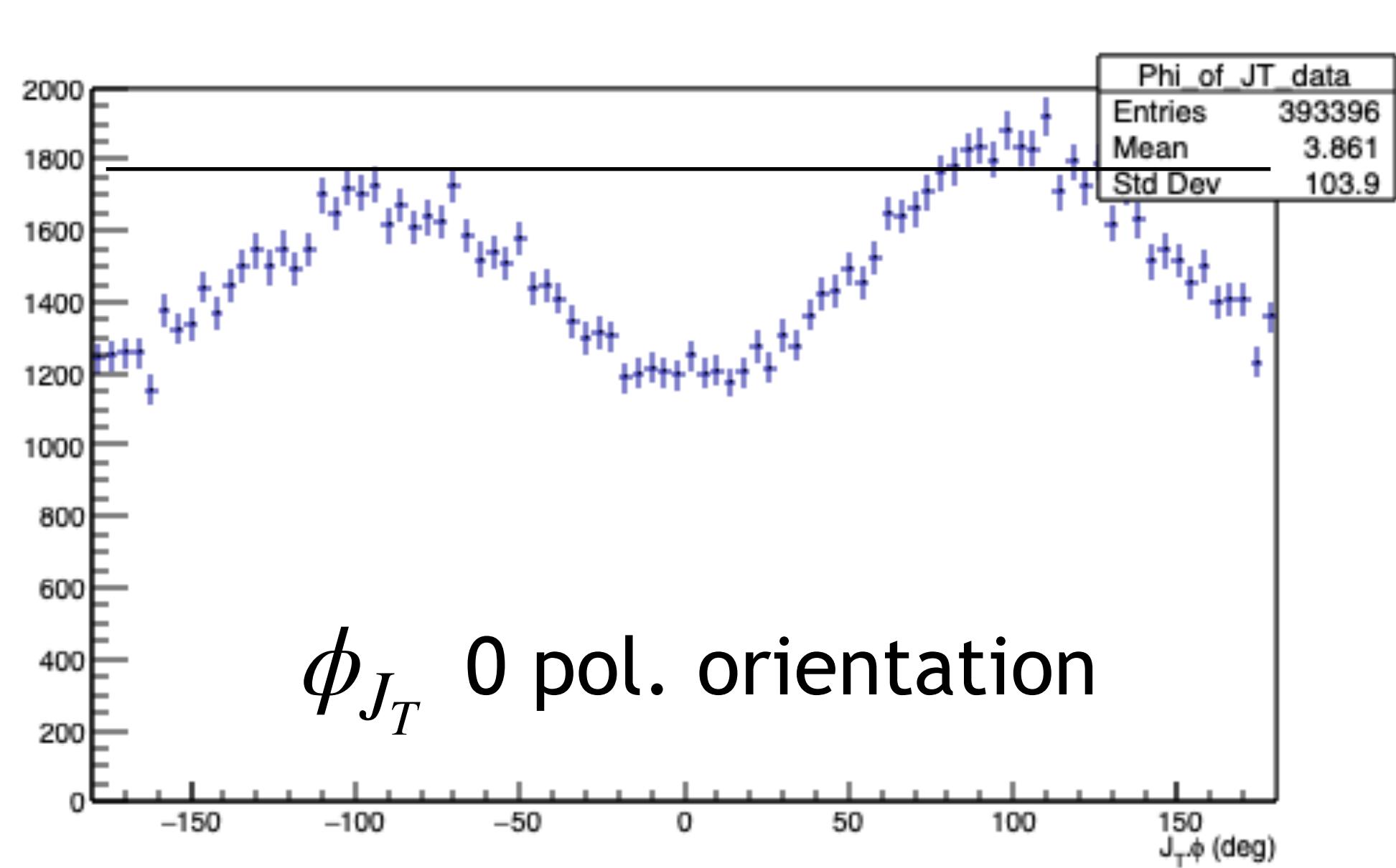
FCAL Elasticity > 0.9

Vertex cut (Window free): $52 < z < 78 \text{ cm}$

e^+e^- Invariant Mass



$\gamma p \rightarrow e^+e^- (p)$ 2018-01 GlueX data, w/ fiducial+N.N. cuts



$$\frac{Y_{\perp}(\phi) - \frac{N_{\perp}}{N_{\parallel}} Y_{\parallel}(\phi)}{Y_{\perp} + \frac{N_{\perp}}{N_{\parallel}} Y_{\parallel}(\phi)} = \frac{\Sigma \cos 2\phi (P_{\perp} + P_{\parallel})}{2 + \Sigma \cos 2\phi (P_{\perp} - P_{\parallel})}$$

$$N_{\perp} = 311346$$

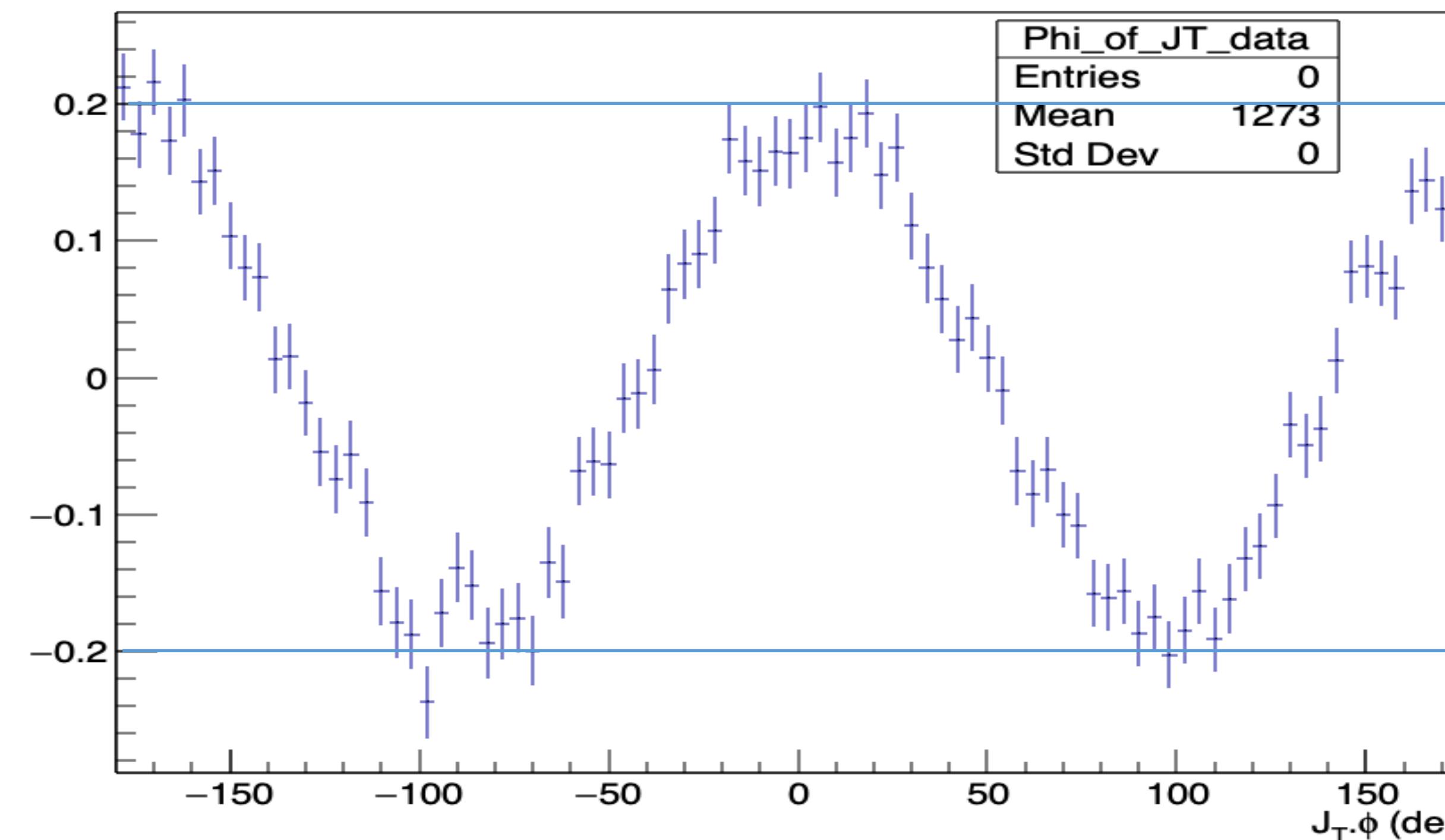
$$N_{\parallel} = 325538$$

$$\frac{N_{\perp}}{N_{\parallel}} = 0.9564$$

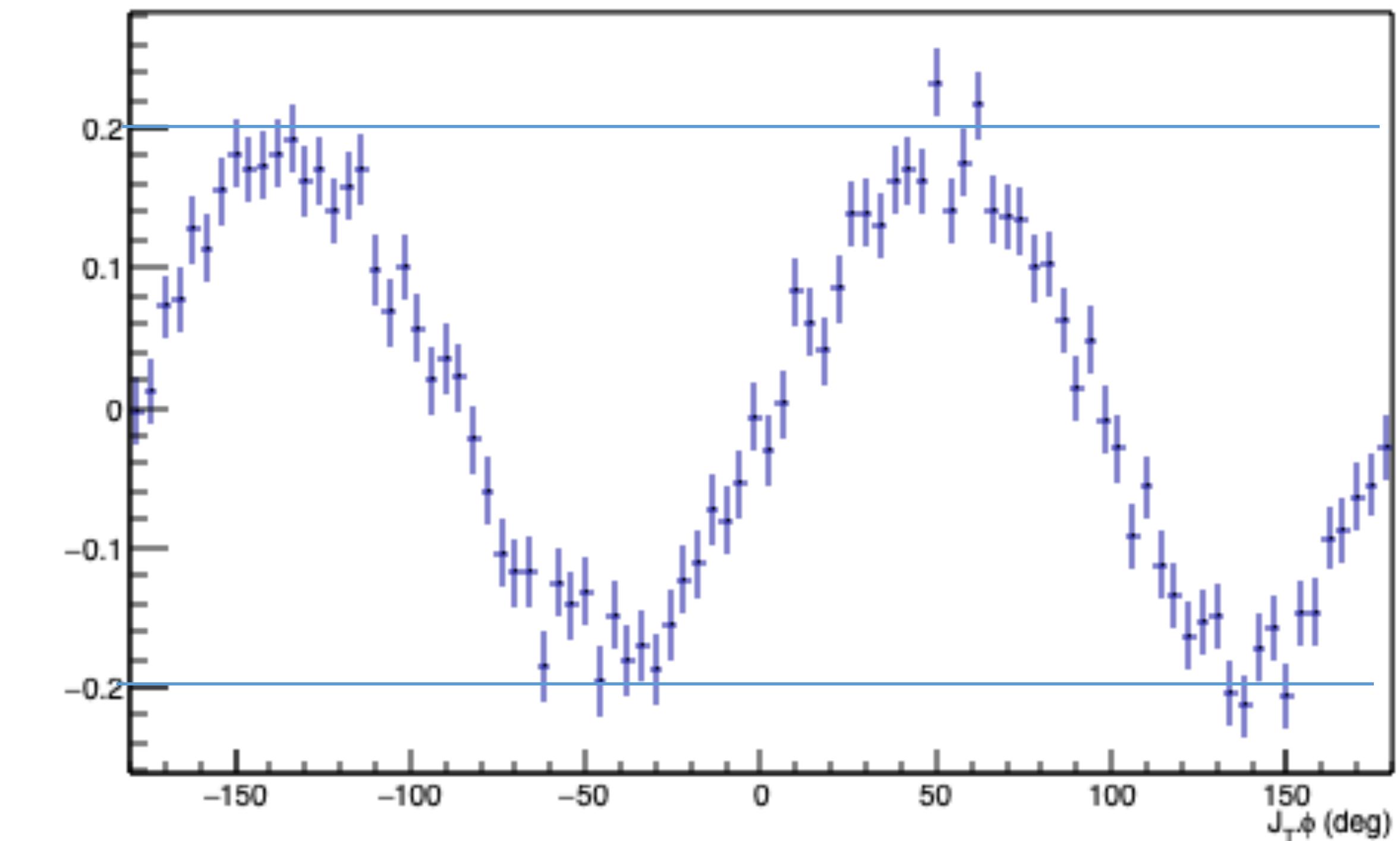
2018-01 GlueX data, $\gamma p \rightarrow e^+e^- (p)$

Yield Asymmetry

0/90 runs

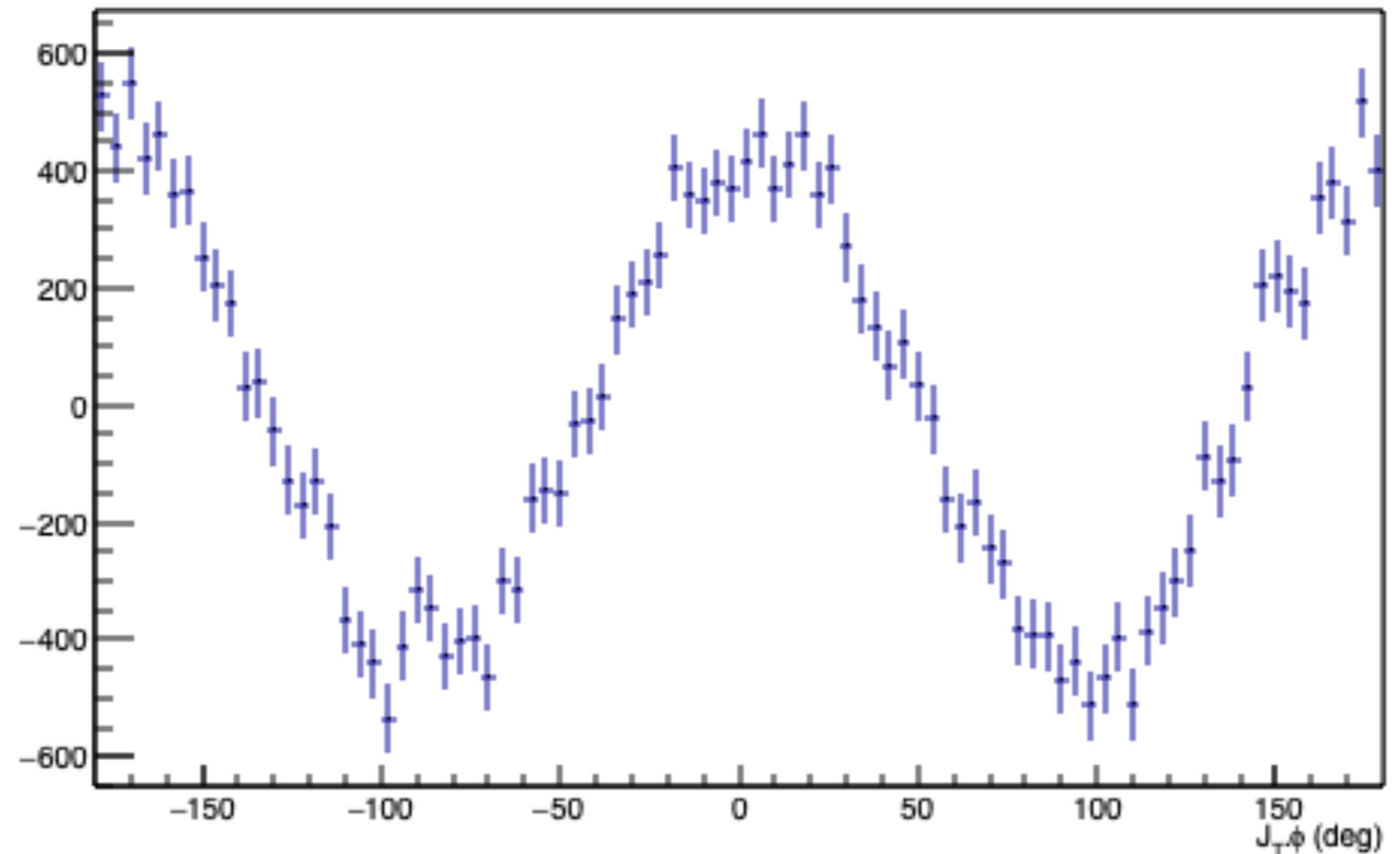


45/135 runs



Just Numerator

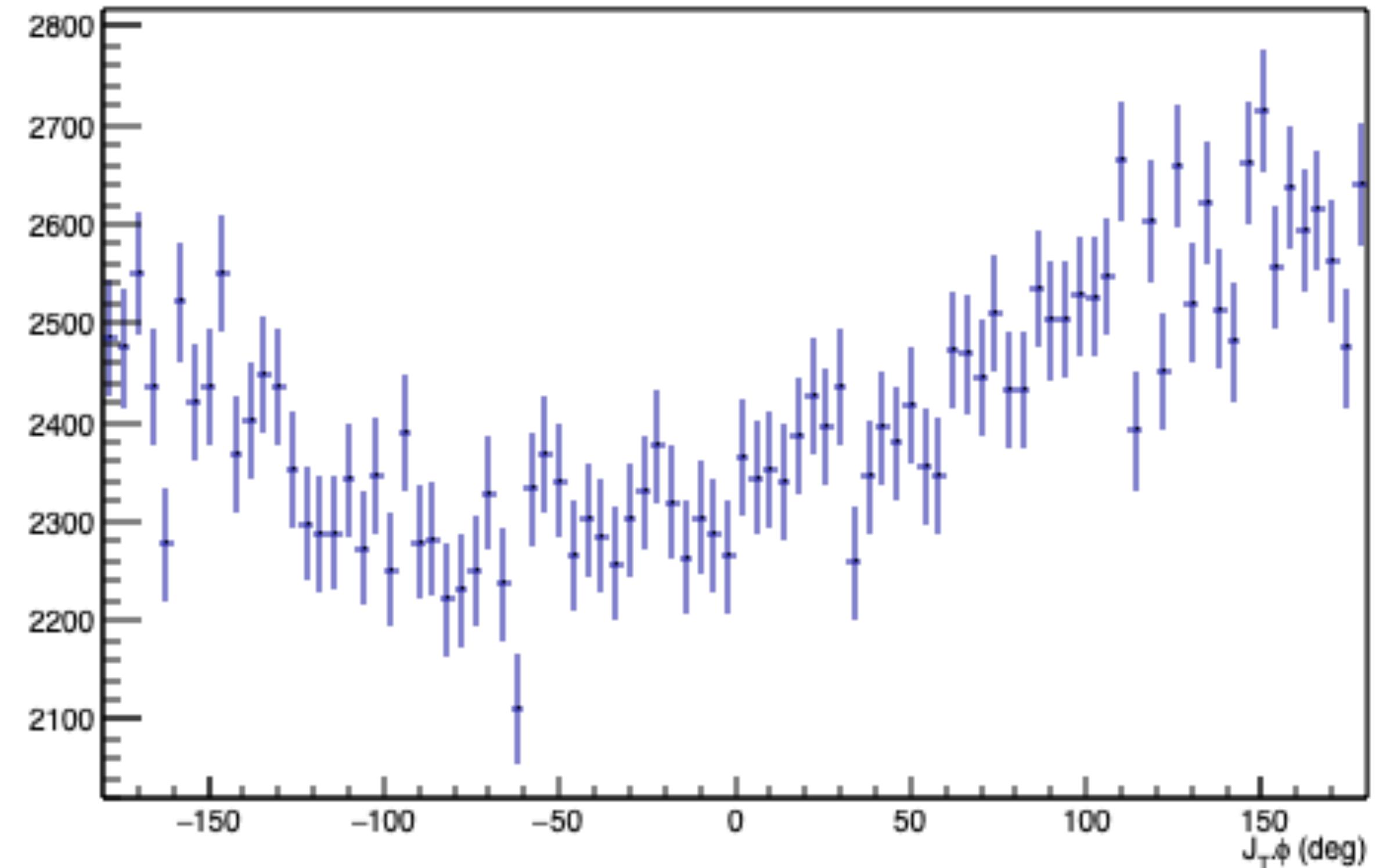
$$Y_{\perp}(\phi) - \frac{N_{\perp}}{N_{\parallel}} Y_{\parallel}(\phi)$$



0/90 pol. Orientation

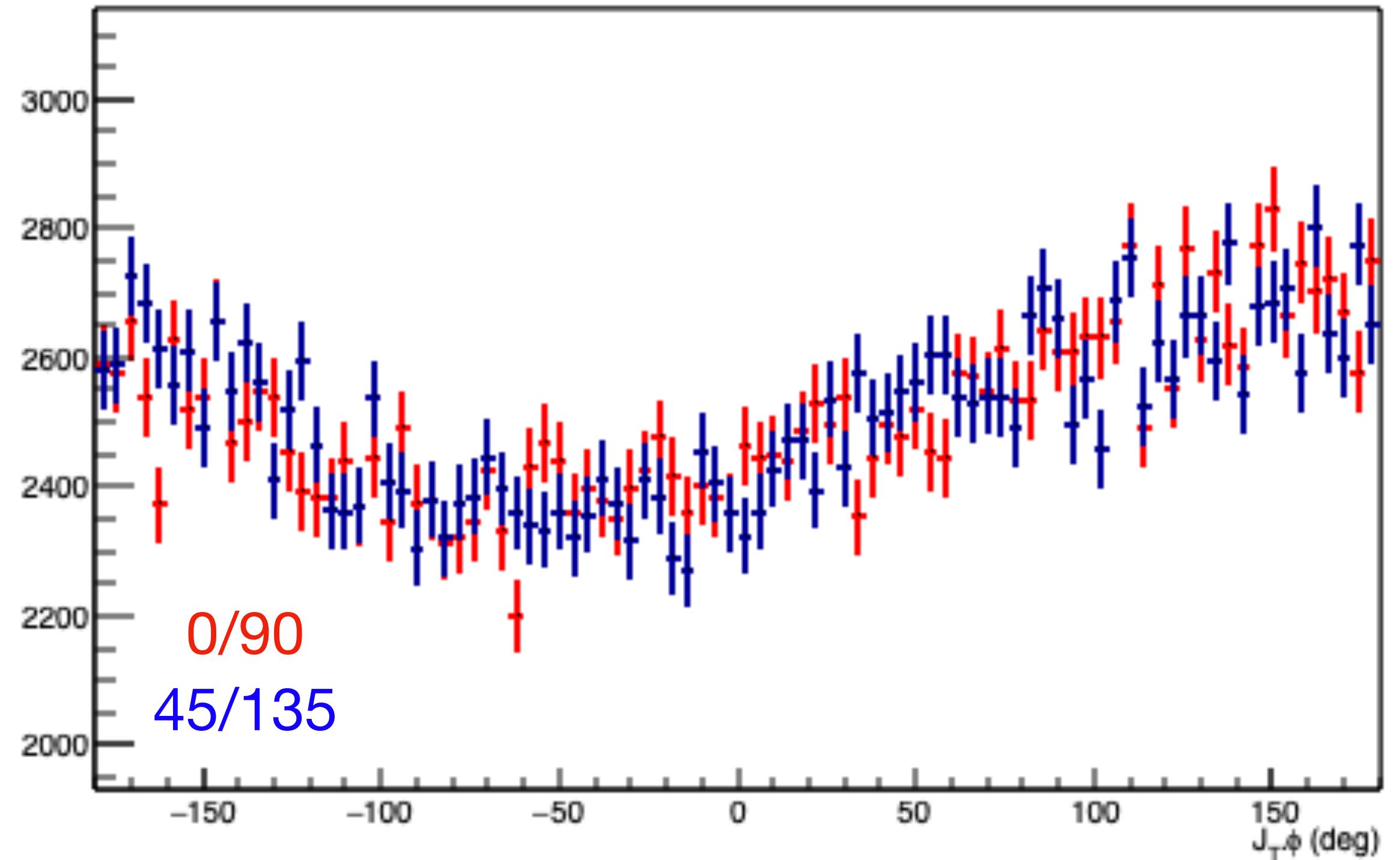
Just Denominator

$$Y_{\perp}(\phi) + \frac{N_{\perp}}{N_{\parallel}} Y_{\parallel}(\phi)$$



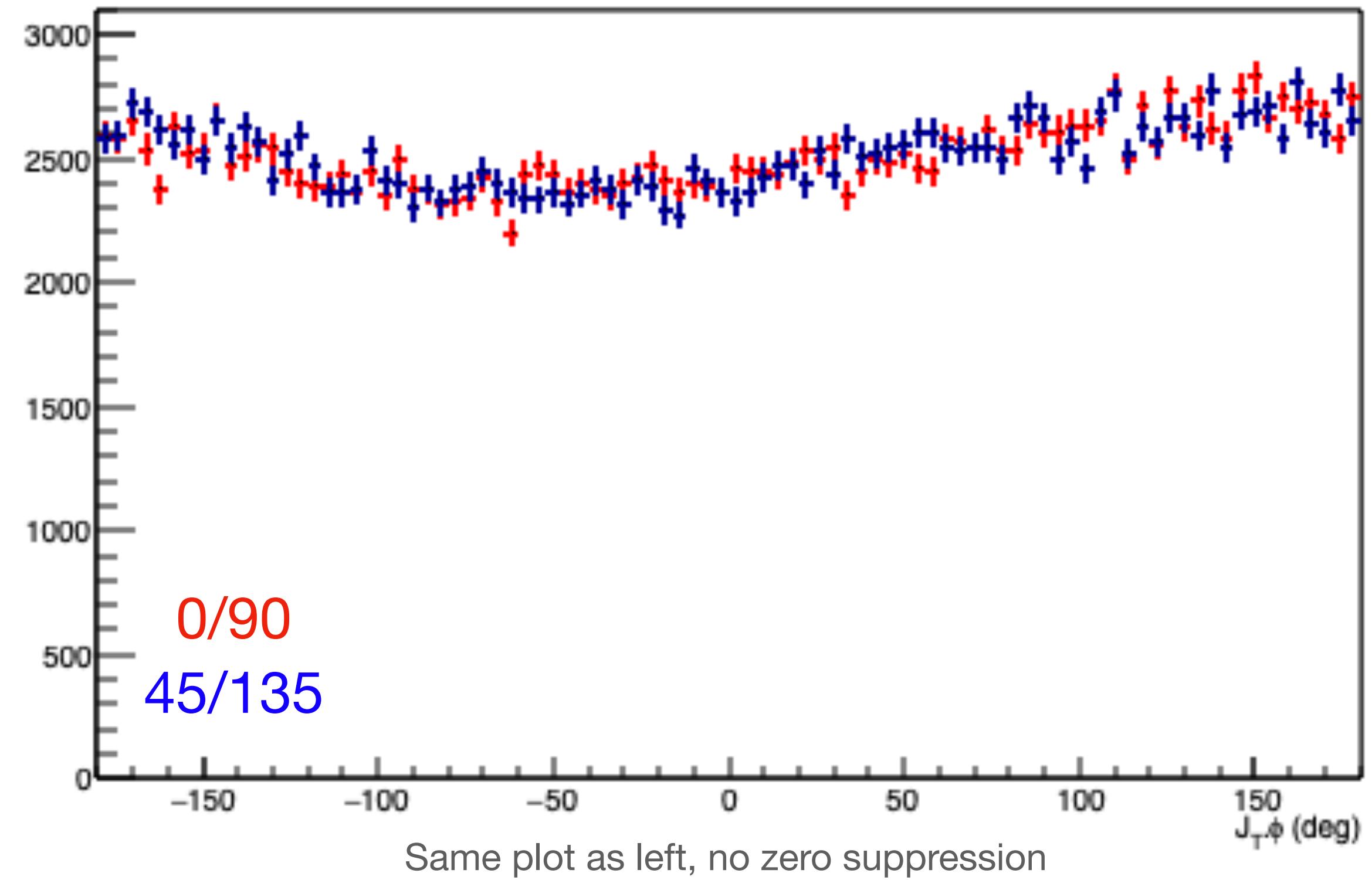
Just Denominator

$$Y_{\perp}(\phi) + \frac{N_{\perp}}{N_{\parallel}} Y_{\parallel}(\phi)$$



0/90

45/135



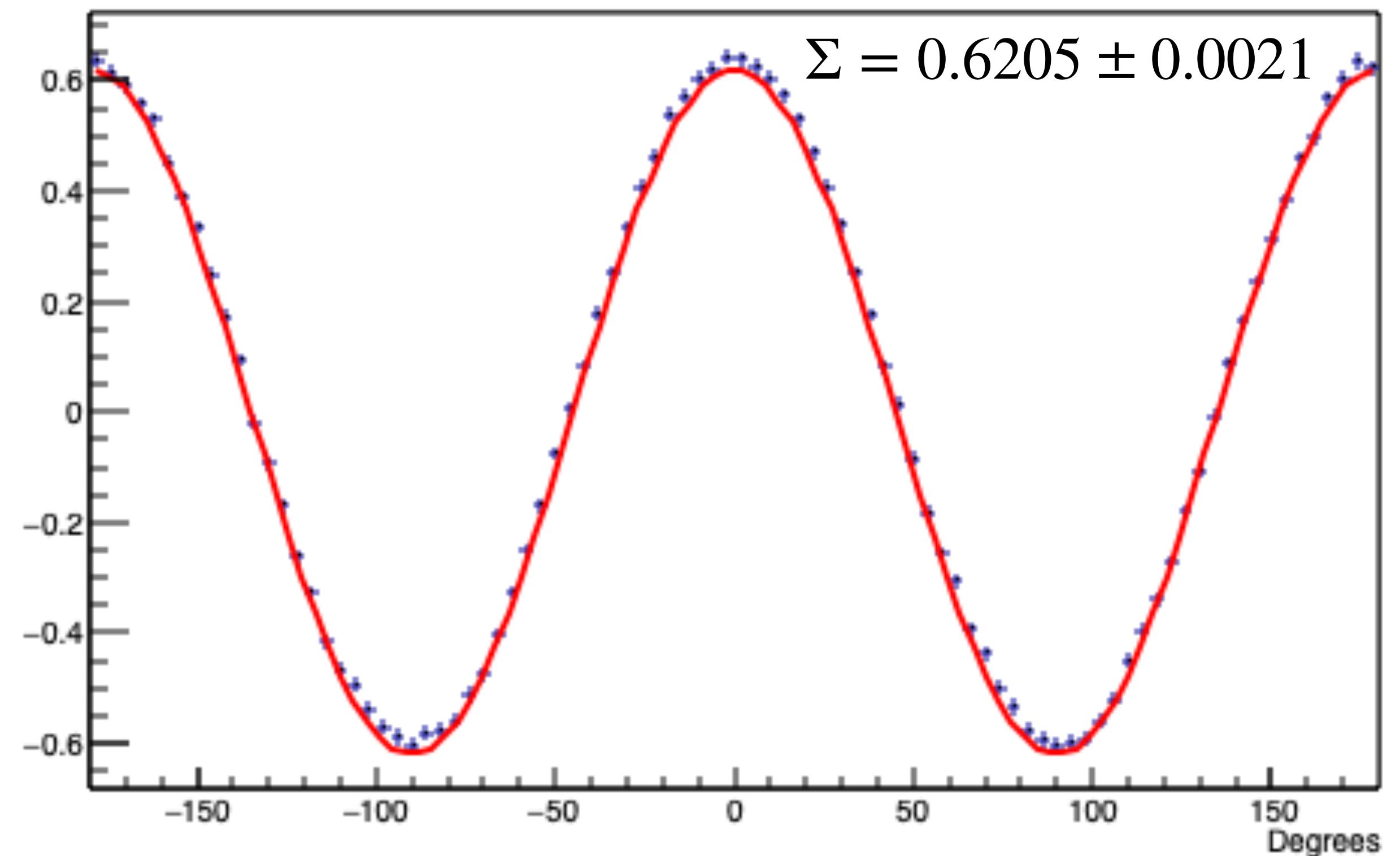
Same plot as left, no zero suppression

$$\frac{Y_{\perp}(\phi) - Y_{\parallel}(\phi)}{Y_{\perp} + Y_{\parallel}(\phi)} = \Sigma \cos 2\phi$$

$$\frac{Y_{\perp}(\phi) - \frac{N_{\perp}}{N_{\parallel}} Y_{\parallel}(\phi)}{Y_{\perp} + \frac{N_{\perp}}{N_{\parallel}} Y_{\parallel}(\phi)} = \frac{\Sigma \cos 2\phi (P_{\perp} + P_{\parallel})}{2 + \Sigma \cos 2\phi (P_{\perp} - P_{\parallel})}$$

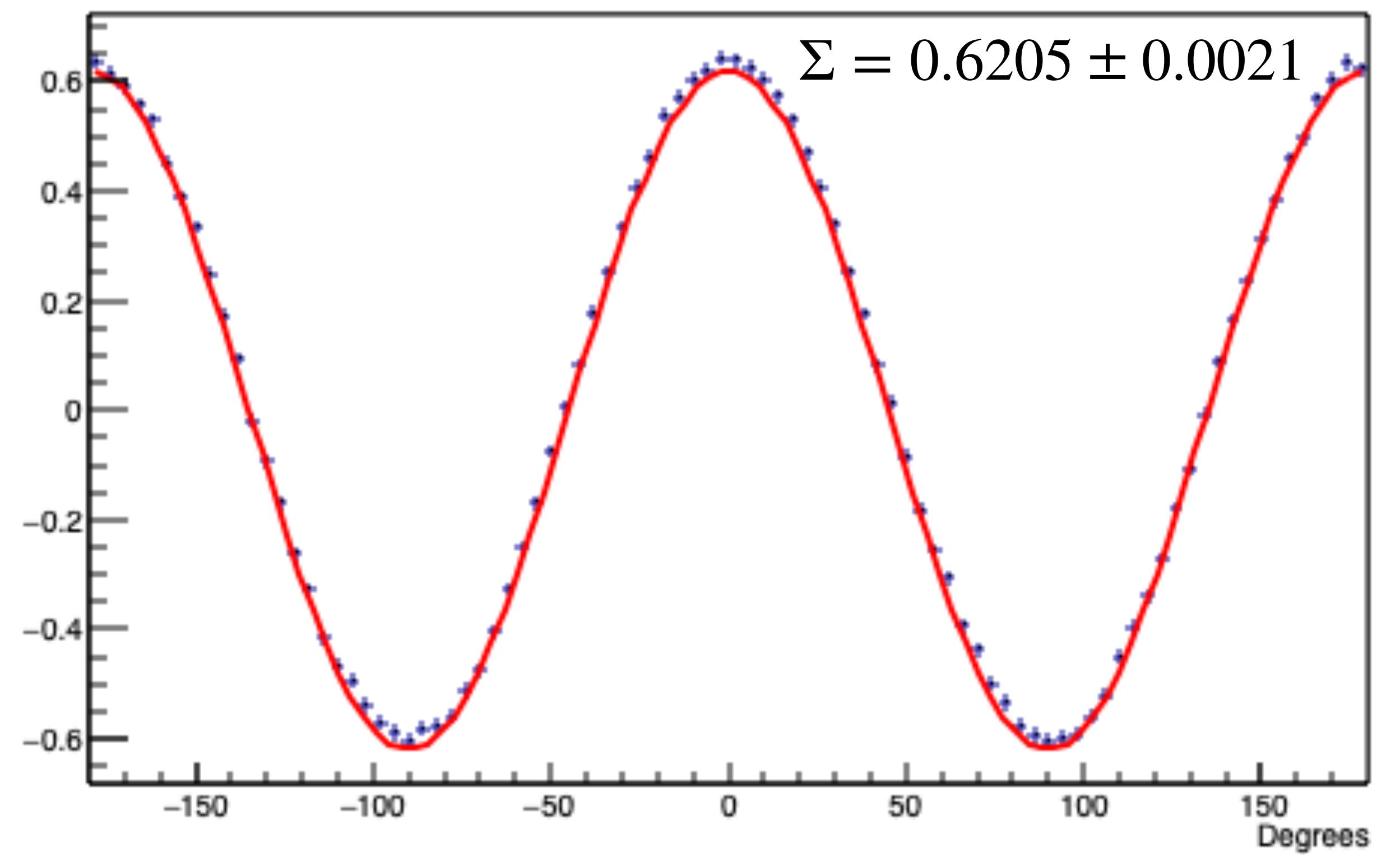
Assume $P_{\perp} - P_{\parallel} = 0$

Simulated Yield Asymmetry



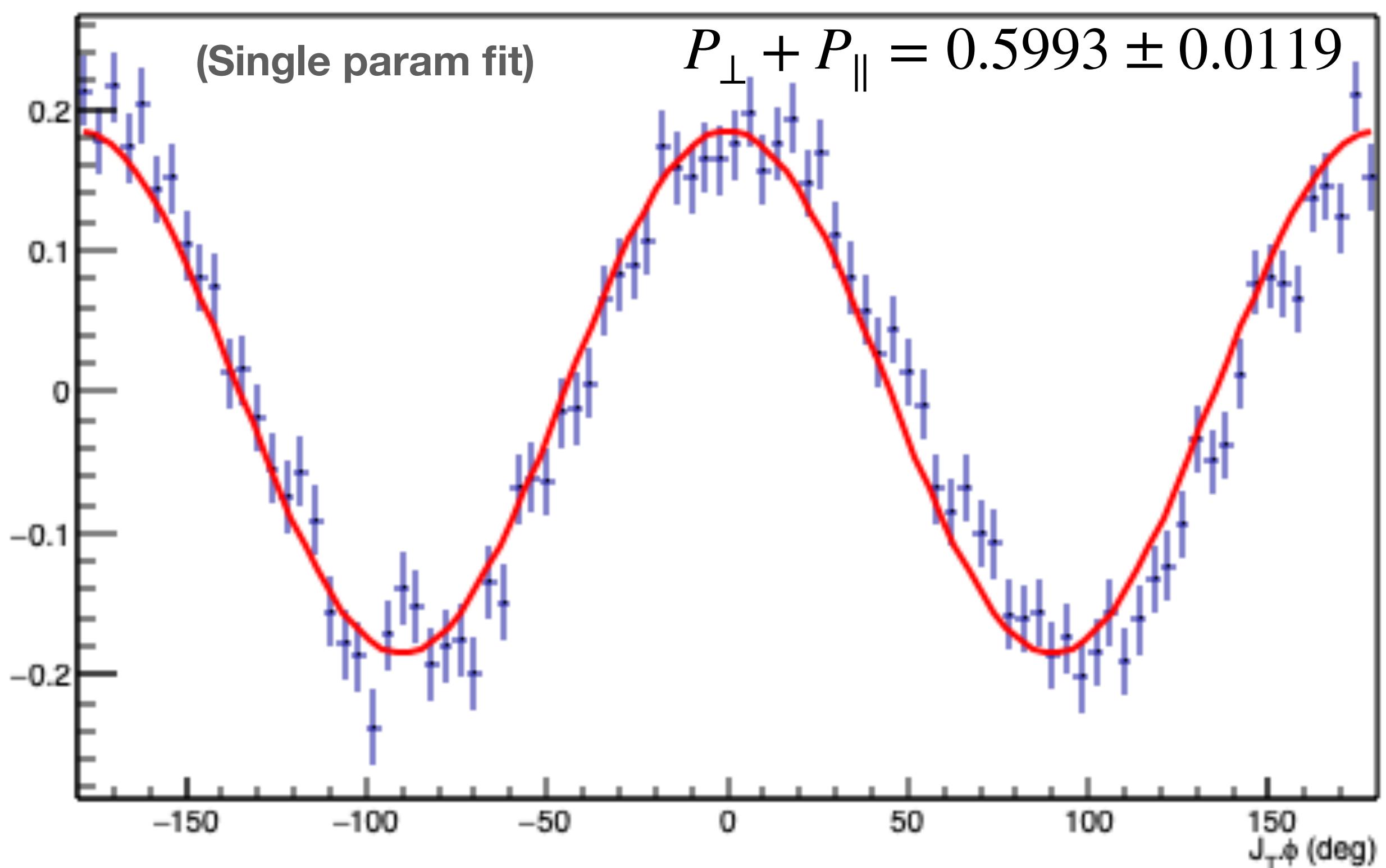
$$\frac{Y_{\perp}(\phi) - Y_{\parallel}(\phi)}{Y_{\perp} + Y_{\parallel}(\phi)} = \Sigma \cos 2\phi$$

Simulated Yield Asymmetry



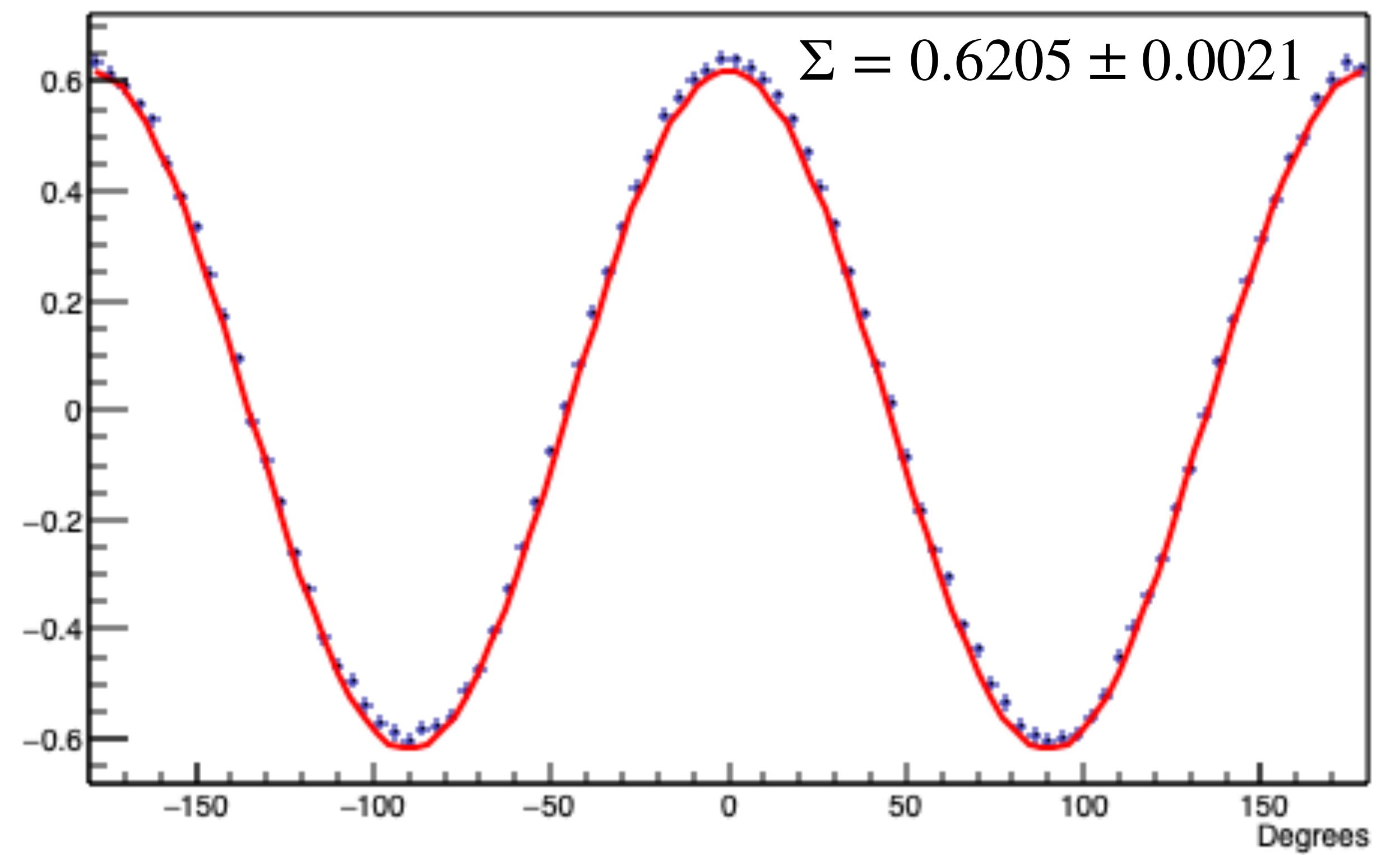
$$\frac{Y_{\perp}(\phi) - \frac{N_{\perp}}{N_{\parallel}} Y_{\parallel}(\phi)}{Y_{\perp} + \frac{N_{\perp}}{N_{\parallel}} Y_{\parallel}(\phi)} = \frac{\Sigma(P_{\perp} + P_{\parallel}) \cos 2\phi}{2}$$

2018-01 Pol = 0 and 90 runs
Data Yield Asymmetry



$$\frac{Y_{\perp}(\phi) - Y_{\parallel}(\phi)}{Y_{\perp} + Y_{\parallel}(\phi)} = \Sigma \cos 2\phi$$

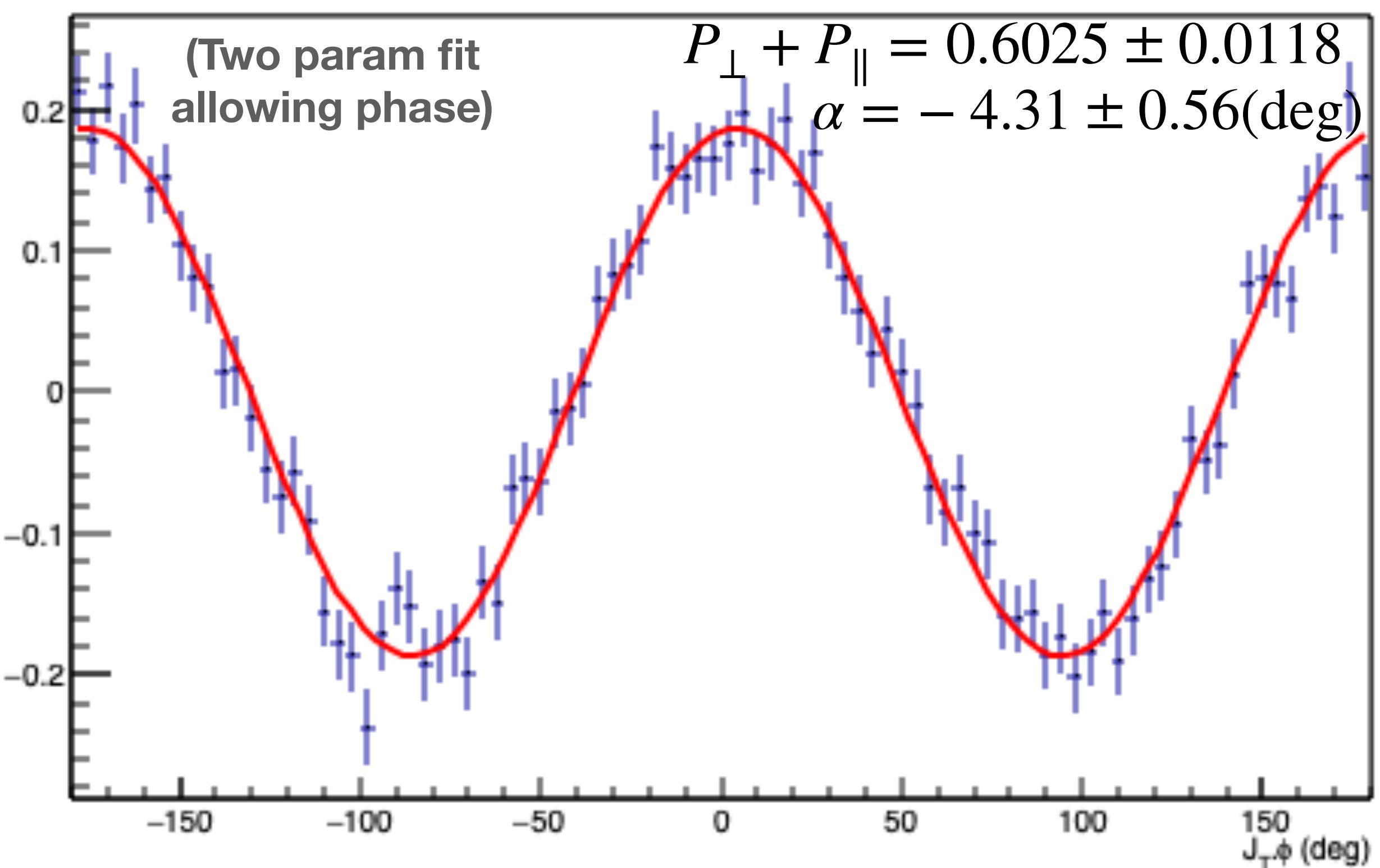
Simulated Yield Asymmetry



$$\frac{Y_{\perp}(\phi) - \frac{N_{\perp}}{N_{\parallel}} Y_{\parallel}(\phi)}{Y_{\perp} + \frac{N_{\perp}}{N_{\parallel}} Y_{\parallel}(\phi)} = \frac{\Sigma(P_{\perp} + P_{\parallel}) \cos 2(\phi + \alpha)}{2}$$

2018-01 Pol = 0 and 90 runs

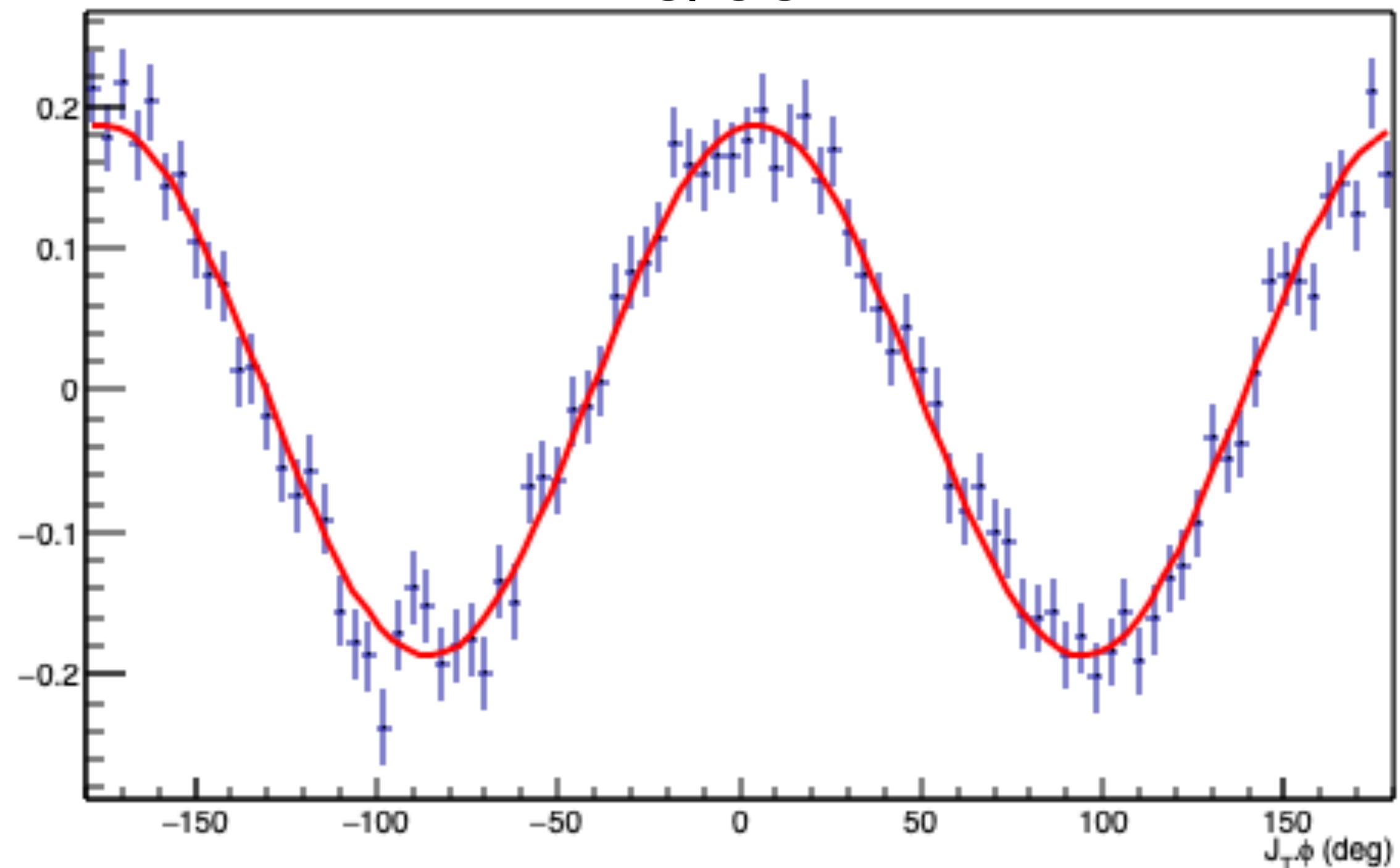
Data Yield Asymmetry



2018-01 GlueX data, $\gamma p \rightarrow e^+e^- (p)$, ϕ_{J_T} Yield Asymmetry

$$\frac{Y_\perp(\phi) - \frac{N_\perp}{N_\parallel} Y_\parallel(\phi)}{Y_\perp + \frac{N_\perp}{N_\parallel} Y_\parallel(\phi)} = \frac{\Sigma(P_\perp + P_\parallel) \cos 2(\phi + \alpha)}{2}$$

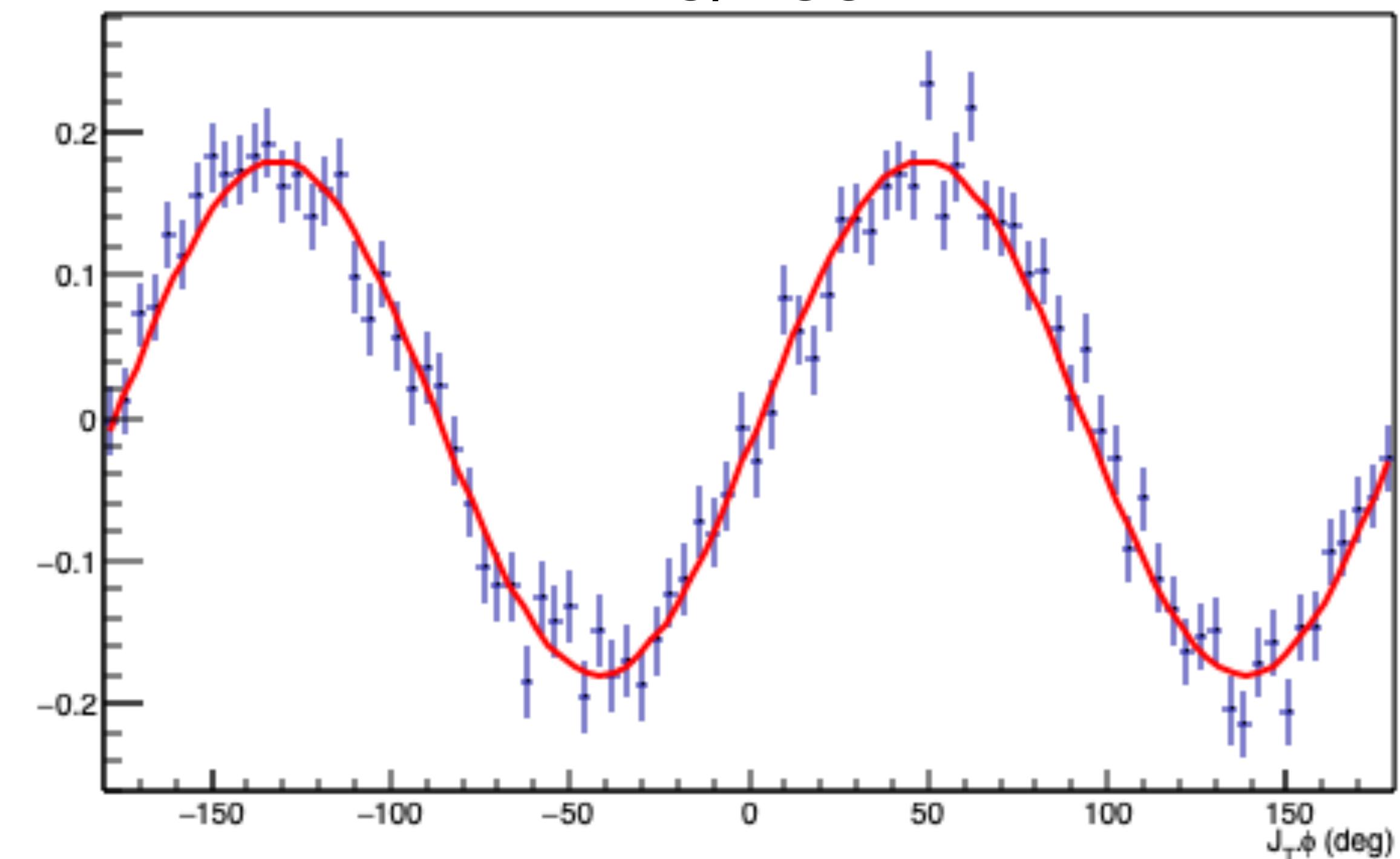
0/90



$$P_\perp + P_\parallel = 0.6025 \pm 0.0118$$

$$\alpha = -4.31 \pm 0.56(\text{deg})$$

45/135



$$P_\perp + P_\parallel = 0.5789 \pm 0.0116$$

$$\alpha = 41.77 \pm 0.57(\text{deg})$$

Preliminary Results

TPOL expected average polarization, 0 and 90 runs

$$\frac{\mathcal{P}_{\perp} + \mathcal{P}_{\parallel}}{2} = 0.341 \pm 0.004$$

BH average polarization; 0 and 90 runs:

$$\frac{\mathcal{P}_{\perp} + \mathcal{P}_{\parallel}}{2} = 0.301 \pm 0.006$$

TPOL expected average polarization, 45 and 135 runs

$$\frac{\mathcal{P}_{\perp} + \mathcal{P}_{\parallel}}{2} = 0.344 \pm 0.004$$

BH average polarization; 0 and 90 runs:

$$\frac{\mathcal{P}_{\perp} + \mathcal{P}_{\parallel}}{2} = 0.290 \pm 0.006$$

Next Steps

Redo study now that improvements in the neural net have been implemented

Rigorously Investigate systematics/stability of result w.r.t. fiducial cuts

Estimated pion contamination is approximately 0.6%. Next, subtract pion yields from the ϕ_{J_T} distributions

Measure analyzing power using Richard's Bethe-Heitler generator to compare with our generator