





# Partial wave analysis studies with simulated $\eta^{(')}\pi^0$ events in GlueX

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# Mesons

# Mesons in standard

quark model



Classified as  $J^{PC}$  multilets:

$$\begin{split} \vec{J} &= \vec{L} + \vec{S} , \\ P &= (-1)^{L+1} \rightarrow \text{Spherical harmonics } (-1)^l \\ &\times \text{Product of individual parites of } q, \overline{q} \ (-1) \\ C &= (-1)^{L+S} \rightarrow \text{Orbital angular momentum } (-1)^l \\ &\times \text{Flip of spin wavefunctions } (-1)^{S+1} \\ &\times \text{ interchanging } q \text{ and } \overline{q} \ (-1) \end{split}$$

J- total angular momentum

S- total quark spin

L- orbital angular momentum between  $q\overline{q}$  pair

P-parity

C - charge conjugation

 $J^{PC}=0^{--}$ , odd<sup>-+</sup> and even<sup>+-</sup> "exotic" quantum numbers are not available.



Quark anti-quark pair coupled to valence gluon. **"Exotic"**  $J^{PC}$  are also available. Predicted by lattice QCD (quantum chromodynamics) calculations (Phys. Rev. D 88, 094505 (2013)).

Primary motivation of the GLUEX is the search for light hybrid mesons.

# $\pi_1$ (1600) results from studies of $\eta'\pi$ system with $\pi$ beam incident on a p target

Evidence for exotic I<sup>G</sup> J<sup>PC</sup> =1<sup>-1-+</sup> state  $\pi_1$  (1600) produced via natural parity exchange (exchanged particle with J<sup>P</sup>s of 0+,1-,2+...) G = C · (-1)<sup>I</sup>, C operator followed by a rotation in isospin (I) E852





(2001)

C. Adolph, et al. [COMPASS Collaboration], Phys. Lett. B740, 303 (2015)

Search for exotic  $\pi_1$  (1600) using the reaction  $\gamma p \rightarrow p \eta' \pi^0$  in GLUEX

The odd waves in  $\eta' \pi^0$  mesonic system have exotic quantum numbers and the lowest of them, the P-wave corresponds to exotic  $\pi_1(1600)$  state.

GLUEX uses linearly polarized photon beam with  $E_{\Upsilon}$ ~ 9GeV



Model for Intensity with polarized photon beam

 $\vec{\gamma}(\lambda, p_{\gamma})p(\lambda_{1}, p_{N}) \rightarrow \pi^{0}(p_{\pi}) \eta(p_{\eta})p(\lambda_{2}, p_{N}')$   $\Phi\text{-angle between } \gamma \text{ polarization vector } \vec{\epsilon}' \text{ and production plane}$   $\Omega\text{- direction of } \eta \text{ in helicity frame}$   $P_{\gamma} \text{ is the degree of linear polarization}$   $A_{\lambda;\lambda_{1}\lambda_{2}}(\Omega)\text{-the reaction amplitude}$   $I(\Omega, \Phi) = \frac{d\sigma}{dtdm_{\eta\pi}d\Omega d\Phi}$   $I(\Omega, \Phi) = I^{0}(\Omega) - P_{\gamma}I^{1}(\Omega)\cos 2\Phi - P_{\gamma}I^{2}(\Omega)\sin 2\Phi$   $I^{0}(\Omega) = \frac{\kappa}{2} \sum_{\lambda,\lambda_{1},\lambda_{2}} A_{\lambda;\lambda_{1}\lambda_{2}}(\Omega) A^{*}_{\lambda;\lambda_{1}\lambda_{2}}(\Omega),$   $I^{1}(\Omega) = \frac{\kappa}{2} \sum_{\lambda,\lambda_{1},\lambda_{2}} A_{-\lambda;\lambda_{1}\lambda_{2}}(\Omega) A^{*}_{\lambda;\lambda_{1}\lambda_{2}}(\Omega),$   $I^{2}(\Omega) = i \frac{\kappa}{2} \sum_{\lambda,\lambda_{1},\lambda_{2}} \lambda A_{-\lambda;\lambda_{1}\lambda_{2}}(\Omega) A^{*}_{\lambda;\lambda_{1}\lambda_{2}}(\Omega),$ 



with  $\kappa$  containing all kinematical factors. The partial wave amplitudes  $T^{l}$  are defined by:

$$A_{\lambda;\lambda_1\lambda_2}(\Omega) = \sum_{lm} T^l_{\lambda m;\lambda_1\lambda_2} Y^m_l(\Omega)$$

We introduce reflectivity basis which allows to trade helicity  $\lambda$  for the reflectivity index  $\epsilon = \pm 1$ , and express helicity amplitudes in terms of reflectivity amplitudes  $T_{-1m;\lambda_1\lambda_2}^l = (-1)^m [{}^{(-)}T_{-m;\lambda_1\lambda_2}^l - {}^{(+)}T_{-m;\lambda_1\lambda_2}^l]$   $T_{+1m;\lambda_1\lambda_2}^l = {}^{(-)}T_{m;\lambda_1\lambda_2}^l + {}^{(+)}T_{m;\lambda_1\lambda_2}^l$ 

At high energies, t-channel exchange and natural (unnatural) exchanges contributes only to the  $\epsilon = +(\epsilon = -)$  components in the reflectivity basis. Define phase rotated spherical harmonics

$$Z_l^m(\Omega, \Phi) \equiv Y_l^m(\Omega) e^{-i\Phi}$$
  
Re $Z_l^m(\Omega, \Phi) = \sqrt{\frac{2l+1}{4\pi}} d_{m0}^l(\theta) \cos(m\varphi - \Phi)$   
Im $Z_l^m(\Omega, \Phi) = \sqrt{\frac{2l+1}{4\pi}} d_{m0}^l(\theta) \sin(m\varphi - \Phi)$ 

Mathieu et al. Phys. Rev. D 100, 054017 (2019)

### Model for Intensity with polarized photon beam Helicity frame Parity invariance implies ${}^{(\epsilon)}T^{l}_{m;-\lambda_{1}-\lambda_{2}} = \epsilon(-1)^{\lambda_{1}-\lambda_{2}} {}^{(\epsilon)}T^{l}_{m;\lambda_{1}\lambda_{2}}$ We take advantage of this constraint to define $l_{m;0}^{(\epsilon)}[l] = {}^{(\epsilon)}T_{m;++}^l \ l_{m;1}^{(\epsilon)}[l] = {}^{(\epsilon)}T_{m;+-}^l$ Are partial wave amplitudes for spin flip k=1 and spin non-flip k=0. $\vec{p}_{\eta}^{CM}$ For each I, there are $2^{2}(2I+1)$ complex partial waves with $\epsilon = \pm 1$ , k=0,1 corresponding $\vec{p}_{\rm recoil}^{CM}$ to target and recoil helicities and m=-l,....l. θ There is no interference between $\epsilon$ =+and $\epsilon$ =- intensities. $\vec{p}_{\text{beam}}^{CM}$ $\vec{p}_{target}^{CM}$ Intensity that involves four coherent sums for each configuration of nucleon spin: ηπ СΜ

$$I(\Omega, \Phi) = 2\kappa \sum_{k} \left\{ \left(1 - P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(-)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 - P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Im}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l}^{m}(\Omega, \Phi)] \right|^{2} + \left(1 + P_{\gamma}\right) \left| \sum_{l,m} [l]_{m;k}^{(+)} \operatorname{Re}[Z_{l$$

Helicity-non-flip amplitudes dominate and we set the helicity-flip amplitudes to zero. This is not restrictive as the target is not polarized in GlueX, and the measured intensities are not sensitive to the details of the nucleon helicity structure.

Natural parity exchanges (corresponding to the amplitudes with  $\epsilon$ =+1) dominate in the energy range of interest.

- Bin data in small bins of  $m_{\eta\pi}$ , t and  $E_{\gamma}$  with constant  $V_{\epsilon LM}$
- Fit data using extended unbinned (in  $(\theta, \varphi)$ ) maximum likelihood method

$$\ln L(V) = \sum_{i=1}^{N} \ln I(V,\theta,\varphi) - \int I(V,\theta,\varphi) \eta(\theta,\varphi) \, d\Omega$$

 $\eta(\theta, \varphi)$  -acceptance

• Minimize –InL using MINUIT, to find V

Mathieu et al. Phys. Rev. D 100, 054017 (2019)

# Generated $2*10^6 (p\eta\pi^0)$ events with AmpTools

## Generated resonances are

- $a_0$ (980 MeV) ٠



# Generated $2*10^6 (p\eta'\pi^0)$ events with AmpTools



For the wave set  $[l]_{m;k}^{(\epsilon)} = \{S_0^{(+)}, P_{0,1}^{(+)}, D_{0,1,2}^{(+)}\}_{k=0}$  with M>= 0

1. Fit intensity to find partial waves using AmpTools.

2. Implement and test calculation of moments in terms of partial waves using the following expressions in terms of the  $\eta' \pi^0$  SDMEs calculated in reflectivity basis: 1 /0

$$H^{0}(LM) = \sum_{\substack{\ell\ell'\\mm'}} \left(\frac{2\ell'+1}{2\ell+1}\right)^{1/2} C^{\ell 0}_{\ell' 0L0} C^{\ell m}_{\ell'm'LM} \rho^{\alpha,\ell\ell'}_{mm'} \qquad \rho^{\alpha,ll'}_{mm'} = \sum_{\epsilon} {}^{(\epsilon)} \rho^{\alpha,ll'}_{mm'} \qquad {}^{(\epsilon)} \rho^{\alpha,\ell\ell'}_{mm'} = \kappa \sum_{k} \left( [\ell]^{(\epsilon)}_{m;k} [\ell']^{(\epsilon)*}_{m;k} + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{-m';k} \right) + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{-m';k} + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{mm'} + (-1)^{m-m'} [\ell]^{(\epsilon)*}_{-m;k} [\ell']^{(\epsilon)*}_{mm'} + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{mm'} + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{mm'} + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{mm'} + (-1)^{m-m'} [\ell]^{(\epsilon)*}_{-m;k} [\ell']^{(\epsilon)*}_{mm'} + (-1)^{m-m'} [\ell]^{(\epsilon)*}_{-m;k} [\ell']^{(\epsilon)*}_{mm'} + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{mm'} + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{-m';k} + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{mm'} + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*$$

3. Compare the results from previous step to moment distributions obtained by Monte Carlo integrations based on the expressions:

$$\begin{split} H^{0}(LM) &= \frac{P_{\gamma}}{2} \int_{\circ} I(\Omega, \Phi) \, d_{M0}^{L}(\theta) \cos M\phi, \\ H^{1}(LM) &= \int_{\circ} I(\Omega, \Phi) \, d_{M0}^{L}(\theta) \cos M\phi \, \cos 2\Phi, \\ \mathrm{Im} \, H^{2}(LM) &= -\int_{\circ} I(\Omega, \Phi) \, d_{M0}^{L}(\theta) \sin M\phi \, \sin 2\Phi, \\ \mathrm{with} \, \int_{\circ} &= (1/\pi P_{\gamma}) \int_{0}^{\pi} \sin \theta \mathrm{d}\theta \int_{0}^{2\pi} \mathrm{d}\phi \int_{0}^{2\pi} \mathrm{d}\Phi \end{split}$$

 $+ (-1)^{m'} [\ell]_{m;k}^{(\epsilon)} [\ell']_{-m';k}^{(\epsilon)*} \Big) ,$  ${}^{(\epsilon)}\rho_{mm'}^{2,\ell\ell'} = -i\epsilon\kappa\sum_{k} \left( (-1)^m [\ell]_{-m;k}^{(\epsilon)} [\ell']_{m';k}^{(\epsilon)*} \right)$  $-(-1)^{m'}[\ell]_{m;k}^{(\epsilon)}[\ell']_{-m';k}^{(\epsilon)*}$ ,  ${}^{(\epsilon)}\rho_{mm'}^{3,\ell\ell'} = \kappa \sum_{k} \left( [\ell]_{m;k}^{(\epsilon)} [\ell']_{m';k}^{(\epsilon)*} \right)$  $-(-1)^{m-m'}[\ell]^{(\epsilon)}_{-m;k}[\ell']^{(\epsilon)*}_{-m';k}$ .

- 4. Compare moments from fitting with true wave set (S0+, P0+, P1+, D0+, D1+, D2+) using good starting values for fit parameters (partial wave amplitudes  $[l]_{m\cdot k}^{(\epsilon)}$ ) to moments from:
  - Fit 1 : fitting with S0+, P0+, D0+, D1+, G0+, G1+ waveset using good starting values for the fit parameters that are common
  - Fit 2: fitting with SO-, PO+, P1+, DO+, D1+, D2+, GO+, G1+ waveset using good starting values for the fit parameters that are 9 common Mathieu et al. Phys. Rev. D 100, 054017 (2019)

# Implementation of calculation of moments

1. Implement and test calculation of moments in terms of partial waves using the following expressions in terms of the  $\eta' \pi^0$ SDMEs calculated in reflectivity basis:

that is applicable for M>=0,  $\epsilon$ >0 and L<=D is coded in "project\_moments\_SPD\_etapi0\_posepsilon". All three codes can be found in halld\_sim/src/programs/AmplitudeAnalysis/

3. I have also added scripts and codes for plotting moments in hd\_utilities/PWA\_scripts/Polarized\_moments\_viaPW

Config file for fitting with generated amplitudes in M and t bins	
define polVal 0.3Typicallyfit FITNAMEreaction EtaPrimePi0 Beam Proton Eta Pi0Can also	refers to unique set of initial and final state particles refer to multiple decay modes of the same set of final state particles
genmc EtaPrimePi0 ROOTDataReader GENMCFILE accmc EtaPrimePi0 ROOTDataReader ACCMCFILE data EtaPrimePi0 ROOTDataReader DATAFILE	Reaction, data reader class, argument Events to fit intensity to
sum EtaPrimePi0 PositiveRe sum EtaPrimePi0 PositiveIm parameter polAngle 1.77 fixed Keywords classes	All amplitudes within a given sum are added coherently
<pre># a0(980) amplitude EtaPrimePi0::PositiveIm::S0+ Zlm 0 0 -1 -1 polAngle polVal amplitude EtaPrimePi0::PositiveRe::S0+ Zlm 0 0 +1 +1 polAngle polVal # a2(1320)a2'(1700) amplitude EtaPrimePi0::PositiveIm::D0+ Zlm 2 0 -1 -1 polAngle polVal amplitude EtaPrimePi0::PositiveRe::D0+ Zlm 2 0 +1 +1 polAngle polVal amplitude EtaPrimePi0::PositiveIm::D1+ Zlm 2 1 -1 -1 polAngle polVal amplitude EtaPrimePi0::PositiveRe::D1+ Zlm 2 1 +1 +1 polAngle polVal amplitude EtaPrimePi0::PositiveRe::D1+ Zlm 2 1 -1 -1 polAngle polVal amplitude EtaPrimePi0::PositiveIm::D2+ Zlm 2 2 -1 -1 polAngle polVal</pre>	Reaction, Sum, amplitude name, amplitude class, arguments Zlm as suggested in GlueX doc-4094 (M. Shepherd) argument 1 : j argument 2 : m argument 3 : real (+1) or imaginary (-1) part
amplitude EtaPrimePi0::PositiveRe::D2+ Zlm 2 2 +1 +1 polAngle polVal # pi1(1600) initialize EtaPrimePi0::PositiveIm::S0+ cartesian 1000.0 0.0 real initialize EtaPrimePi0::PositiveRe::S0+ cartesian 1000.0 0.0 real	argument 4 : 1 + (+1/-1) * P_gamma argument 5 : polarization angle (in Deg.) argument 6 : beam properties config file or fixed polarization
initialize EtaPrimePi0::PositiveIm::D0+ cartesian 70.0 70.0 initialize EtaPrimePi0::PositiveRe::D0+ cartesian 70.0 70.0 initialize EtaPrimePi0::PositiveIm::D1+ cartesian 70.0 70.0	— Initial value of partial wave amplitudes $[l]_{m;k}^{(\epsilon)}$ in cartesian coordinate system
initialize EtaPrimePi0::PositiveRe::D1+ cartesian 70.0 70.0 initialize EtaPrimePi0::PositiveIm::D2+ cartesian 70.0 70.0 initialize EtaPrimePi0::PositiveRe::D2+ cartesian 70.0 70.0	Factors with the same reaction sum and amplitude name are multiplied together
constrain EtaPrimePi0::PositiveIm::S0+ EtaPrimePi0::PositiveRe::S0+	Same amplitudes corresponding to different sums should be equal

• • •

# Results for bin M=1.37 and t<0.3

Amplitudes used in fitting are S0+, P0+, P1+, D0+, D1+, D2+. Good starting values for fit parameters Fit results



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Fit 1 results (fitting in M and t bins)



Analysis strategy

For the wave set  $[l]_{m;k}^{(\epsilon)} = \left\{ S_0^{(+)}, P_{0,1}^{(+)}, D_{0,1,2}^{(+)} \right\}_{k=0}$  with M>= 0

1. Fit intensity to find partial waves using AmpTools.

2. Implement and test calculation of moments in terms of partial waves using the following expressions in terms of the  $\eta' \pi^0$  SDMEs calculated in reflectivity basis:

$$H^{0}(LM) = \sum_{\substack{\ell\ell'\\mm'}} \left(\frac{2\ell'+1}{2\ell+1}\right)^{1/2} C^{\ell 0}_{\ell' 0L0} C^{\ell m}_{\ell'm'LM} \rho^{\alpha,\ell\ell'}_{mm'} \qquad \rho^{\alpha,\ell\ell'}_{mm'} = \sum_{\epsilon} {}^{(\epsilon)} \rho^{\alpha,ll'}_{mm'} \qquad {}^{(\epsilon)} \rho^{\alpha,\ell\ell'}_{mm'} = \kappa \sum_{k} \left( [\ell]^{(\epsilon)}_{m;k} [\ell']^{(\epsilon)*}_{m';k} + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{-m';k} + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{-m';k} + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{mm'} + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m} [\ell']^{(\epsilon)*}_{mm'} + (-1)^{m-m'} [\ell']^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)}_{-m;k} + (-1)^{m-m'} [\ell']^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)}_{-m;k} + (-1)^{m-m'} [\ell']^{(\epsilon)}_{-m;k} + (-1)^{($$

3. Compare the results from previous step to moment distributions obtained by Monte Carlo integrations based on the expressions:

$$\begin{split} H^{0}(LM) &= \frac{P_{\gamma}}{2} \int_{\circ} I(\Omega, \Phi) \, d_{M0}^{L}(\theta) \cos M\phi, \\ H^{1}(LM) &= \int_{\circ} I(\Omega, \Phi) \, d_{M0}^{L}(\theta) \cos M\phi \, \cos 2\Phi, \\ \mathrm{Im} \, H^{2}(LM) &= -\int_{\circ} I(\Omega, \Phi) \, d_{M0}^{L}(\theta) \sin M\phi \, \sin 2\Phi, \\ \mathrm{with} \, \int_{\circ} &= (1/\pi P_{\gamma}) \int_{0}^{\pi} \sin \theta \mathrm{d}\theta \int_{0}^{2\pi} \mathrm{d}\phi \int_{0}^{2\pi} \mathrm{d}\Phi \end{split}$$

$$\begin{aligned} & \left| \left( -1 \right)^{m-m'} [\ell]_{-m;k}^{(\epsilon)} [\ell']_{-m';k}^{(\epsilon)*} \right) , \\ & \left| \left( \epsilon \right) \rho_{mm'}^{1,\ell\ell'} = -\epsilon \kappa \sum_{k} \left( (-1)^{m} [\ell]_{-m;k}^{(\epsilon)} [\ell']_{m';k}^{(\epsilon)*} \right) \\ & + (-1)^{m'} [\ell]_{m;k}^{(\epsilon)} [\ell']_{-m';k}^{(\epsilon)*} \right) , \\ & \left( \epsilon \right) \rho_{mm'}^{2,\ell\ell'} = -i\epsilon \kappa \sum_{k} \left( (-1)^{m} [\ell]_{-m;k}^{(\epsilon)} [\ell']_{m';k}^{(\epsilon)*} \right) \\ & - (-1)^{m'} [\ell]_{m;k}^{(\epsilon)} [\ell']_{-m';k}^{(\epsilon)*} \right) , \\ & \left( \epsilon \right) \rho_{mm'}^{3,\ell\ell'} = \kappa \sum_{k} \left( [\ell]_{m;k}^{(\epsilon)} [\ell']_{m';k}^{(\epsilon)*} \\ & - (-1)^{m-m'} [\ell]_{-m;k}^{(\epsilon)} [\ell']_{-m';k}^{(\epsilon)*} \right) , \end{aligned}$$

4. Compare moments from fitting with true wave set (S0+, P0+, P1+, D0+, D1+, D2+) using good starting values for fit parameters (partial wave amplitudes) to moments from:

- Fit 1 : fitting with S0+, P0+, D0+, D1+, G0+, G1+ waveset using good starting values for the fit parameters that are common
- Fit 2: fitting with S0-, P0+, P1+, D0+, D1+, D2+, G0+, G1+ waveset using good starting values for the fit parameters that are common
   Mathieu et al. Phys. Rev. D 100, 054017 (2019)

Polarized moments calculated with partial waves







Polarized moments calculated with partial waves

Analysis strategy

For the wave set  $[l]_{m;k}^{(\epsilon)} = \left\{ S_0^{(+)}, P_{0,1}^{(+)}, D_{0,1,2}^{(+)} \right\}_{k=0}$  with M>= 0

1. Fit intensity to find partial waves using AmpTools.

2. Implement and test calculation of moments in terms of partial waves using the following expressions in terms of the  $\eta' \pi^0$  SDMEs calculated in reflectivity basis:

$$H^{0}(LM) = \sum_{\substack{\ell\ell' \\ mm'}} \left( \frac{2\ell'+1}{2\ell+1} \right)^{1/2} C^{\ell 0}_{\ell' 0L0} C^{\ell m}_{\ell'm'LM} \rho^{\alpha,\ell\ell'}_{mm'} \qquad \rho^{\alpha,\ell\ell'}_{mm'} = \sum_{\epsilon} {}^{(\epsilon)} \rho^{\alpha,ll'}_{mm'} \qquad {}^{(\epsilon)} \rho^{\alpha,\ell\ell'}_{mm'} = \kappa \sum_{k} \left( [\ell]^{(\epsilon)}_{m;k} [\ell']^{(\epsilon)*}_{m';k} + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{-m';k} \right) + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{mm'} \qquad {}^{(\epsilon)} \rho^{1,\ell\ell'}_{mm'} = -\epsilon \kappa \sum_{k} \left( (-1)^{m} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{m';k} \right) + (-1)^{m-m'} [\ell]^{(\epsilon)}_{-m;k} [\ell']^{(\epsilon)*}_{m';k} \right)$$

3. Compare the results from previous step to moment distributions obtained by Monte Carlo integrations based on the expressions:

$$\begin{split} H^{0}(LM) &= \frac{P_{\gamma}}{2} \int_{\circ} I(\Omega, \Phi) \, d_{M0}^{L}(\theta) \cos M\phi, \\ H^{1}(LM) &= \int_{\circ} I(\Omega, \Phi) \, d_{M0}^{L}(\theta) \cos M\phi \, \cos 2\Phi, \\ \operatorname{Im} H^{2}(LM) &= -\int_{\circ} I(\Omega, \Phi) \, d_{M0}^{L}(\theta) \sin M\phi \, \sin 2\Phi, \\ \operatorname{with} \, \int_{\circ} &= (1/\pi P_{\gamma}) \int_{0}^{\pi} \sin \theta \mathrm{d}\theta \int_{0}^{2\pi} \mathrm{d}\phi \int_{0}^{2\pi} \mathrm{d}\Phi \end{split}$$

$$\frac{1}{k} \left( (-1)^{m-m'} [\ell]_{-m;k}^{(\epsilon)} [\ell']_{-m';k}^{(\epsilon)*} \right) + (-1)^{m-m'} [\ell]_{-m;k}^{(\epsilon)} [\ell']_{-m';k}^{(\epsilon)*} + (-1)^{m'} [\ell]_{-m;k}^{(\epsilon)} [\ell']_{m';k}^{(\epsilon)*} + (-1)^{m'} [\ell]_{m;k}^{(\epsilon)} [\ell']_{-m';k}^{(\epsilon)*} \right) ,$$

$$\stackrel{(\epsilon)}{=} \rho_{mm'}^{2,\ell\ell'} = -i\epsilon\kappa \sum_{k} \left( (-1)^{m} [\ell]_{-m;k}^{(\epsilon)} [\ell']_{m';k}^{(\epsilon)*} - (-1)^{m'} [\ell]_{m;k}^{(\epsilon)} [\ell']_{-m';k}^{(\epsilon)*} \right) ,$$

$$\stackrel{(\epsilon)}{=} \rho_{mm'}^{3,\ell\ell'} = \kappa \sum_{k} \left( [\ell]_{m;k}^{(\epsilon)} [\ell']_{m';k}^{(\epsilon)*} - (-1)^{m-m'} [\ell]_{-m;k}^{(\epsilon)} [\ell']_{-m';k}^{(\epsilon)*} \right) ,$$

4. Compare moments from fitting with true wave set (S0+, P0+, P1+, D0+, D1+, D2+) using good starting values for fit parameters (partial wave amplitudes) to moments from:

- Fit 1 : fitting with S0+, P0+, D0+, D1+, G0+, G1+ waveset using good starting values for the fit parameters that are common
- Fit 2: fitting with S0-, P0+, P1+, D0+, D1+, D2+, G0+, G1+ waveset using good starting values for the fit parameters that are common
   Mathieu et al. Phys. Rev. D 100, 054017 (2019)

Fitting data with S0-, P0+, P1+, D0+, D1+, D2+ amplitude set with S0-, P0+, D0+, D1+, G0+, G1+.

Fit 2 results (fitting in M and t bins)



# Bootstrapping method for estimation of uncertainties



- 1. Draw a Bootstrap Sample from the original sample data with replacement with size n.
- 2. Evaluate intensity for each Bootstrap Sample which will result in B estimates of intensity.
- 3. Construct a histogram of B estimates of intensity and use it to make further statistical inference, such as:

• Estimating the standard error of statistic for Intensity.

Distributions of moment values from 100 bootstrapping samples for M bin=5 and t bin=1



Polarized moments calculated with partial waves with uncertainties from bootstrapping



Polarized moments calculated with partial waves with uncertainties from bootstrapping

0 < t < 0.3 (GeV/c)<sup>2</sup> Calculated from fitted amplitudes , with bootstrapping uncert.

Obtained by weighting events of data sample





Fitting data with S0-, P0+, P1+, D0+, D1+, D2+ amplitude set with S0-, P0+, P1+, D0+, D1+, D2+, G0+, G1+.

# Fit 1 results (fitting in M and t bins)

Amplitudes used in fitting are S0+, P0+, P1+, D0+, D1+, D2+. Good starting values for fit parameters



Distributions of moment values from 100 bootstrapping samples for M bin=4 and t bin=1



Polarized moments calculated with partial waves with uncertainties from bootstrapping



Polarized moments calculated with partial waves with uncertainties from bootstrapping

0 < t < 0.3 (GeV/c)<sup>2</sup> Calculated from fitted amplitudes , with bootstrapping uncert.

Obtained by weighting events of data





